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ENDF Formats and Procedures for Photon Production and Interaction Data


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# ENDF Formats and Procedures for Photon Production and Interaction Data 

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#### Abstract

The ENDF formats for storage of photon production and photon interaction data are presented as are procedures for entering these data. Formats for photon production data are a revision of those given in Los Alamos Scientific Laboratory Report LA-3801 (ENDF-111) and supersede the earlier format. Sample sets of photon production data are given in the new format. Procedures are given to minimize ambiguities in entering data in the format and to facilitate writing processing codes to use the data. Classes of data for which formats and procedures are specified include photon production cross sections, multiplicities, transition probability arrays, angular distributions, energy spectra, and energy-angle distributions; also, photon interaction cross sections, secondary angular, energy, energy-angle distributions, and atomic form factors are included.


1. INTRODUCTION

Since the genesis of a Shielding Subcommittee of the Cross-Section Evaluation Working Group (CSEWG) in early 1967, the question of formats for storing photon production and photon interaction data has received considerable attention. A format emphasizing multiplicities (yields) and transition probabilities was adopted and documented in detail. ${ }^{1}$ However, in early 1969, the Subcommittee decided to recommend a revised format, isomorphic to the existing one, but placing photon production cross sections, multiplicities, and energy distributions in separate files. This report, the second of two volumes, documents the latest accepted format and provides recommended procedures for placing data in the format and for writing codes to process the data. General features and descriptions of the data formats and record types are given in Chaps. 1 through 5 of the companion (BNL 50274, ENDF-102, Revised, Vol. I) to this report. A knowledge of those chapters is presumed for this volume.

The format for photon interaction data remains essentially unchanged, but it is included in this report for both completeness and clarification of the procedures.

The arrangement of this report is similar to that proposed to the CSEWG for the revision of the ENDF neutron data format specifications. ${ }^{2}$ Sections 2 and 3 contain specifications for photon production and photon interaction files, respectively. Discussion of each ENDF photon data file is subdivided into two subsections, one for formats and one for procedures. A glossary of symbols and variable names is provided in Appendix A for convenient reference. The variable names are, in most cases, consistent with the corresponding FORTRAN variables in the photon production data processing codes. 3

The same units as specified by Honeck ${ }^{2}$ are used where appropriate, and the set is expanded as follows:

| energy | -eV |
| :--- | :--- |
| angle | - dimensionless cosine |
| cross section | - barns |
| mass | - units of the neutron mass |
| yield | - photons |

In any file, sections are always arranged in increasing order of MT number. Structures of sections are given below in the form of the binary records. Unused fields are now denoted by "b" (blank) in the format prescriptions, as opposed to the integer and floating-point zeros ( 0 or 0.0 ) used previously for ENDF/B. The

BCD card-image formats for each section are given in Appendix $B$, of ten with a sample set of data in the format (see Vol. I, Sec. 5.3 for a detailed description of the card-image formats).
2. PHOTON PRODUCTION

Photon production data are divided into five distinct files.

File

12 Multiplicities and transition probability arrays
13 Photon production cross sections
14 Photon angular distributions
15 Continuous photon energy spectra
16 Photon energy-angle distributions

With the exception of File 12, all of the files are closely analogous to the corresponding neutron data files with the same number (modulo 10). The purpose of File 12 is to provide additional methods for representing the energy dependence of photon production cross sections. The parallel structure between neutron and photon data files facilitates the use of existing neutron data retrieval routines ${ }^{4}$ for the photon data. Record types are identical to those used for neutron data (e.g., HEAD, CONT, LIST, TAB1, TAB2). The allowed reaction type (MT) numbers are the same as those assigned for Files 1 through 7, and they continue to maintain as closely as possible the parallelism with the United Kingdom Atomic Energy Authority data file. 5 However, they may have somewhat different meanings for photon production that require additional explanation in some cases, as follows
(1) $M T=3$ should be used in Files 12 through 16 to represent composite ( $n, x y$ ) cross sections, that is, photon production cross sections from more than one reaction type that have been lumped together.
(2) There is no apparent reason to have redundant or derived data for the photon production files as is the case for the neutron files, i.e., MT = 3, 4, etc. Therefore, to avoid confusion, the join of all sections of Files 12 and 13 should represent the photon production, with each section being disjoint with all others.
(3) Let us consider how one might represent the inelastic gamma-ray production data. The differential cross section for producing a gama ray of energy $E_{\gamma}$ resulting from the excitation of the $m_{o}$ th level of the residual nucleus and the subsequent transition between two definite levels ( $j \rightarrow i$ ), including the effects of cascading from the $m_{0}-j$ levels higher than $j$, is

$$
\begin{equation*}
\frac{d \sigma}{d E \gamma}\left(E_{\gamma}, E, m_{o}, i, j\right)=\delta\left(E_{\gamma}-E_{j}+\varepsilon_{i}\right) A_{j, i} \sigma_{m_{o}}(E) \overbrace{\ell=1}^{m_{o-j}} \sum_{m_{\ell}=j}^{m_{\ell-1}-1} T P_{m_{\ell-1}}, m_{\ell} \tag{1}
\end{equation*}
$$

where

$$
\begin{aligned}
& \sigma_{m_{0}}(E) \equiv \text { neutron cross sections for exciting the } m_{o} \text { th level with } \\
& \text { neutron energy } E \text {, } \\
& \delta\left(E_{\gamma}-\varepsilon_{j}+\varepsilon_{i}\right) \equiv \underset{\substack{\text { delta function with } \\
\text { residual nucleus, }}}{ }, \varepsilon_{i} \text { being energy levels of the } \\
& \begin{aligned}
& \mathrm{TP} \\
& \mathrm{k}, \ell \equiv \text { probability of the residual nucleus having a transition } \\
& \text { to the } \ell \text { th level given that it was initially in the ex- }
\end{aligned} \\
& \text { cited state corresponding to the kth level, and } \\
& A_{k, \ell} \equiv \text { probability of a gamma ray of energy } E_{\gamma}=\varepsilon_{k}-\varepsilon_{\ell} \text { being } \\
& \text { emitted as a result of the residual nucleus having a } \\
& \text { transition from the kth to the } \ell \text { th level. }
\end{aligned}
$$

We are at once beset with no clear choice of ENDF representation in terms of section number. The data may naturally be identified both with the $m_{o}$ th level and the $j$ th level. To avoid this problem, we can sum Eq. (1) over $m_{o}$
$\frac{d \sigma_{1}}{d_{\gamma}}\left(E_{\gamma}, E, i, j\right)=\sum_{m_{o}=j}^{N} \frac{d \sigma}{d E_{\gamma}}\left(E_{\gamma}, E, m_{o}, i, j\right)$,
where $N$ is the highest level that can be excited by a neutron of incident energy $E$ (i.e., $\varepsilon_{N} \leq \frac{A W R}{A W R+1} E$ ). This gives us a de-excitation cross section that has the characteristic of singling out a definite gamma-ray transition and has the advantage when experimental data are to be represented. The de-excitation cross section is identified with the $j$ th level. Alternatively, we can sum Eq. (1) over $i$ and $j$
$\frac{d \sigma_{2}}{d E_{\gamma}}\left(E_{\gamma}, E, m_{o}\right)=\sum_{j=1}^{m_{o}} \sum_{i=0}^{j-1} \frac{d \sigma}{d E_{\gamma}}\left(E_{\gamma}, E, m_{o}, i, j\right) \quad$.
This gives us an excitation cross section that has the characteristic of singling out a definite excited state and has the advantage when calculated data are to be represented. The excitation cross section is identified with the $m_{0}$ th level. If Eq. (2) is summed over $i$ and $j$, or equivalently Eq. (3) is summed over $m_{o}$, then

$$
\begin{equation*}
\frac{d \sigma}{d E_{\gamma}}\left(E_{\gamma}, E\right)=\sum_{m_{0}=1}^{N} \frac{d \sigma_{2}}{d E_{\gamma}}\left(E_{\gamma}, E, m_{o}\right) \equiv \sum_{j=1}^{N} \sum_{i=0}^{j-1} \frac{d \sigma_{1}}{d E_{\gamma}}\left(E_{\gamma}, E, i, j\right) . \tag{4}
\end{equation*}
$$

This gives us a cross section for all possible excitations and transitions and thus corresponds to the total inelastic neutron cross section for discrete levels.

It is recommended that $M T=4$ be used for the data represented by Eq. (4), as well as for the continuum. If, however, it is expedient or useful to use $\mathrm{MT}=51$ through 91 , then one must use either the de-excitation cross sections of Eq. (2) or the excitation cross sections of Eq. (3) but not both. A restriction occurs if the transition probability array option is used and if the entire neutron energy range is not covered by the known transition probabilities. Then, for $M T=51$ through 90 in File 12 to be used for the remaining neutron energy range, a representation by excitation multiplicities must be used.

The integrated cross sections of File 13 are obtained by integrating Eqs. (1) through (4) over $\mathrm{E}_{\boldsymbol{\gamma}}$.
(4) The same remarks as given in Item (3) apply for discrete gamma rays from ( $n, p \gamma$ ), ( $n, d \gamma$ ), ( $n, t \gamma$ ), ( $n, 3 \mathrm{He} \mathrm{\gamma}$ ), and ( $n, \alpha \gamma$ ) reactions, and the use of $\mathrm{MT}=103,104,105,106$, and 107 is recommended for these cases.

### 2.1. File 12: Multiplicities and Transition Probability Arrays

File 12 can be used to represent the neutron energy dependence of photon production cross sections by means of either multiplicities or transition probability arrays. Both methods rely upon the use by processing codes of neutron cross sections from File 2 and/or File 3 in order to generate absolute photon production cross sections.

Multiplicities can be used to represent the cross sections of discrete photons and/or the integrated cross sections of continuous photon spectra. The MT numbers that appear in File 12 designate the particular neutron cross sections (File 2 and/or File 3) to which the multiplicities are referred. The use of multiplicites is the recommended method of presenting ( $n, \gamma$ ) capture gamma-ray cross sections provided, of course, that the ( $n, \gamma$ ) cross section is adequately represented in File 2 and/or File 3.

For well-established level decay schemes, the use of transition probability arrays offers a concise method for presenting ( $n, x y$ ) information. With this method, the actual decay scheme of the residual nucleus for a particular reaction (defined by MT number) is entered into File 12. This information can then be used by a processing code together with discrete level excitation cross sections from File 3 to calculate discrete gamma-ray production cross sections. This option cannot be used to represent the integrals of continuous photon spectra.

### 2.1.1. File 12 Format

File 12 is made up of sections with each section giving information for a particular reaction type (MT number). The information in each section is given either as multiplicities ( $L \varnothing=1$ ) or as transition probability arrays ( $L \emptyset=2$ ). Each section always starts with a HEAD record and ends with a SEND record.

### 2.1.1.1. Option 1 ( $L \varnothing=1$ ): Multiplicities

The neutron energy dependence of photon production cross sections is represented by tabulating a set of neutron energy and multiplicity pairs $\left[E, y_{k}(E)\right]$ for each discrete photon and for the photon energy continuum.* The subscript $k$ designates a particular discrete photon or a photon continuum, and the total number of such sets is represented by NK.

The multiplicity or yield $y_{k}(E)$ is defined by

$$
y_{k}(E)=\frac{\sigma_{k}^{\gamma}(E)}{\sigma(E)} \quad \text { (photons), }
$$

where $E$ designates neutron energy and $\sigma(E)$ is the neutron cross section in File 2 and/or File 3 to which the multiplicity is referred (by the MT number). In the case of discrete photons, $\sigma_{k}^{\gamma}(E)$ is the photon production cross section for the discrete photon designated by $k$. In the case of photon continua, $\sigma_{k}^{\gamma}(E)$ is the cross section for the photon continuum integrated over photon energy. In the continuum case,

$$
\begin{aligned}
y_{k}(E) & =\frac{\sigma_{k}^{\gamma}(E)}{\sigma(E)}=\frac{\int \frac{d \sigma_{k}^{\gamma}}{d E_{\gamma}}\left(E_{\gamma}+E\right) d E_{\gamma}}{\sigma(E)} \\
& =\frac{\int \sigma(E) y_{k}\left(E_{\gamma}+E\right) d E_{\gamma}}{\sigma(E)}=\int_{0}^{E_{\gamma}^{\max }} y_{k}\left(E_{\gamma}+E\right) d E_{\gamma},
\end{aligned}
$$

[^0]where $E_{\gamma}$ designates photon energy $(e V), \frac{d \sigma_{k}^{\gamma}}{d E_{\gamma}}\left(E_{\gamma} \leftarrow E\right)$ is the absolute photon energy distribution in barns/eV, and $y_{k}\left(E_{\gamma} \leftarrow E\right)$ is the relative energy distribution in photons/eV. The quantity $y_{k}\left(E_{\gamma} \leftarrow E\right)$ can be broken down further as
$$
y_{k}\left(E_{\gamma} \leftarrow E\right)=y_{k}(E) f_{k}\left(E_{\gamma} \leftarrow E\right),
$$
which results in the requirement that
$$
\int_{0}^{E_{\gamma}^{\max }} f_{k}\left(E_{\gamma}+E\right) d E_{\gamma}=1
$$

Any time a continuum representation is used for a given MT number in either File 12 or 13 , then the normalized energy distribution $f_{k}\left(E_{\gamma} \leftarrow E\right)$ must be given in File 15 under the same MT number.

As a check quantity, the total yield

$$
Y(E)=\sum_{k=1}^{N K} y_{k}(E) \quad \text { (photons) }
$$

is also tabulated for each MT number if $\mathrm{NK}>1$.
The structure of a section for $L \emptyset=1$ is

and the structure of each subsection is

$$
\left[M A T, 12, M T / E G_{k}, E S_{k} ; b, L F ; N R, N P / E_{i n t} / y_{k}(E)\right] T A B 1,
$$

[^1]where

| $k$ level is unknown or if a continuous photon spectrum is produced, $\mathrm{ES}_{k} \equiv 0.0$ should be used. <br> $E G_{k} \equiv$ the photon energy. For a continuous photon energy distribution, $E G_{k} \equiv 0.0$ should be used. <br> LF $\equiv$ the photon energy distribution law number, which presently has onl two values definedı $\begin{aligned} \mathrm{LF} & =1, \text { a normalized tabulated function (in File 15), and } \\ & =2, \text { a discrete photon energy. } \end{aligned}$ |
| :---: |
|  |  |
|  |  |
|  |  |

### 2.1.1.2. Option $2(L \emptyset=2)$ : Transition Probability Arrays

With this option, the only data required are the level energies, deexcitation transition probabilities, and (where necessary) conditional photon emission probabilities. Given this information, the photon energies and their multiplicities can readily be calculated. Also, photon production cross sections can then be computed for any given level from the excitation cross section in File 3, along with the transition probability array. Similarly, multiplicities and photon production cross sections can be constructed for the total cascade. For any given level, the transition and photon emission probability data given in the section are for photons originating at that level only; any further cascading is determined from the data for the lower levels.

Now define the following variables.
$L G=1$, simple case (all transitions are $\gamma$ emission).
$=2$, complex case (internal conversion or other competing processes occur).
NS 三 number of levels below the present one, including the ground state. (The present level is also uniquely defined by the MT number and by its energy level.)
NT $\equiv$ number of transitions for which data are given in a list to follow (i.e., number of nonzero transition probabilities), NT $\leq$ NS.
$E S_{i} \equiv$ energy of the $i t h$ level, $i=0(1) N S . \quad\left(E S_{0} \equiv 0.0\right.$, the ground state.)
$\mathrm{TP}_{\mathrm{i}} \equiv \mathrm{TP}_{\mathrm{i}, \mathrm{NS}, \mathrm{I}}, \mathrm{the}$ probability of a direct transition from level NS to level
$G P_{i} \equiv G P_{N S}, i$, the probability that, given a transition from level NS to level $i$, the transition is a photon transition (i.e., the conditional probability of photon emission).

$$
A_{i} \equiv\left(T P_{i}\right)\left(G P_{i}\right)
$$

Note that each level can be identified by its NS number. Then the energy of a photon from a transition to level $i$ is given by $E_{\gamma}=E S_{N S}-E S_{i}$, and its multiplicity is given by $y\left(E_{\gamma} \leftarrow E\right)=\left(T P_{i}\right)\left(G P_{i}\right)$. It is implicitly assumed that the transition probability array is independent of incident neutron energy.

The structure of a section for $L \varnothing=2$ is

$$
\left.\left\{\begin{array}{l}
{[\mathrm{MAT}, 12, \mathrm{MT} / \mathrm{ZA}, \mathrm{AWR} ; \mathrm{L} \varnothing=2, \mathrm{LG} ;}
\end{array} \mathrm{NS}, \mathrm{~b}\right] H E A D .\right] .
$$

If $L G=1$, the array $B_{i}$ consists of NT doublets ( $\left.E S_{i}, T_{i}\right)$; if $L G=2$, it consists of NT triplets ( $E S_{i}, \mathrm{TP}_{i}, \mathrm{GP}_{\mathrm{i}}$ ). Here the subscript $i$ is a running index over the levels below the level for which the transition probability array is being given (i.e., below level NS). The doublets (ES ${ }_{i}, T_{i}$ ) or triplets $\left(E S_{i}, \mathrm{TP}_{i}, \mathrm{GP}_{\mathrm{i}}\right)$ are given in decreasing magnitude of energy $E S_{i}$.

### 2.1.2. File 12 Procedures

1. Under Option 1 , the subsections are given in decreasing magnitude of $E G_{k}$.
2. Under Option 1 , the convention is that the subsection for the continuum photons, if present, is last. In this case, the last value of $E G_{k}\left(E G_{N K}\right)$ is set equal to 0.0 , and logical consistency with Procedure 1 is maintained.
3. Under Option 1 , the values of $E G_{k}$ should be consistent to within four significant figures with the corresponding $E G_{k}$ values for the File 14 photon angular distributions. This allows processing and "physics" checking codes to match photon yields with the corresponding angular distributions.
4. Under Option $1, E S_{k}$ is the energy of the level from which the photon originates. If $E S_{k}$ is unknown or not meaningful (as for the continuous photon spectrum), the value 0.0 should be entered.
5. If capture and fission resonance parameters are given in File 2, photon production for these reactions should be given by use of Option 1 of File 12, instead of as photon production cross sections in File 13. This is due to the voluminous data required to represent the resonance structure in File 13 and the difficulty of calculating multigroup photon production matrices from such data.
6. Under Option 1 , the total yield table, $Y(E)$, should exactly span the same energy range as the combined energy range of all the $y_{k}(E)$. Within that range,

$$
Y(E)=\sum_{k=1}^{N K} y_{k}(E)
$$

should hold within four significant figures.
7. The excitation cross sections for all the levels appearing in the transition probability arrays must, of course, be given in File 3.
8. The join of all sections, regardless of the option used, should represent the photon production data with no redundancy. For example, $M T=4$ cannot include any photons given elsewhere under $\mathrm{MT}=51$ through 91. Likewise, there can be no redundancy between Files 12 and 13.
9. If only one energy distribution is given under Option $1(N K=1)$, then the TAB1 record for the $Y(E)$ table is simply deleted to avoid repetitive entries.
10. Data should not be given in File 12 for reaction types that do not appear in Files 2 and/or 3.
11. Under Option 2, the level energies, $E S_{i}$, in the transition probability arrays are given in decreasing magnitude.
12. The MT numbers for which transition probability data are given should be for consecutive levels, beginning at the first level, with no embedded levels omitted.
13. The energies of photons arising from level transitions should be consistent within four significant figures with the corresponding $E G_{k}$ values in File 14. Therefore, care must be taken to specify level energies to the appropriate number of significant figures.
14. Under Option 2, the sum of the transition probabilities ( $\mathrm{TP}_{i}$ ) over i should equal 1.0000 (that is, should be unity to within five significant figures).
15. The limit on the number of energy points in any table of $Y(E)$ or $y_{k}(E)$ is 1000. This is an upper limit that should rarely be approached in practice because yields are normally smoothly varying functions of incident neutron energy.
16. The limit on the number of interpolation regions is 10 .
17. Tabulations of nonthreshold data should normally cover at least the energy range $10^{-5} \mathrm{eV} \leq \mathrm{E} \leq 2 \times 10^{7} \mathrm{eV}$, where practical. Threshold data should be given from threshold energy up to $2 \times 10^{7} \mathrm{eV}$, where practical.

### 2.2 File 13: Photon Production Cross Sections

The purpose of File 13 is the same as that for File 12; namely, it can be used to represent the neutron and photon energy dependence of photon production cross sections. In File 13, however, absolute cross sections in barns are tabulated, and there is no need to refer back to the neutron files.

### 2.2.1. File 13 Format

Similar to File 12, File 13 is made up of sections with each section giving information for a particular reaction type (MT number). Each section always starts with a HEAD record and ends with a SEND record.

The representation of the energy dependence of the cross sections is accomplished by tabulating a set of neutron energy-cross section pairs [ $\left.E, \sigma_{k}^{\gamma}(E)\right]$ for each discrete photon and for the photon energy continuum. The subscript $k$ designates particular discrete photons or the photon continuum, and the total number of such sets is NK. In the case of discrete photons, $\sigma_{k}^{\gamma}(E)$ is the photon production cross section (b) for the photon designated by $k$. In the case of the photon continuum, $\sigma_{k}^{\gamma}(E)$ is the integrated (over photon energy) cross section for the photon continuum* designated by $k$. In the continuum case,

$$
\begin{equation*}
\sigma_{k}^{\gamma}(E)=\int_{0}^{E_{\gamma}^{\max }} \frac{d \sigma_{k}^{\gamma}}{d E_{\gamma}}\left(E_{\gamma}+E\right) d E_{\gamma} \tag{b}
\end{equation*}
$$

where $E_{\gamma}$ designates photon energy ( $e V$ ), and $\frac{d \sigma_{k}^{\gamma}}{d E_{\gamma}}\left(E_{\gamma} \leftarrow E\right)$ is the absolute photon energy distribution in b/eV. The energy distribution can be further broken down as

$$
\frac{d \sigma_{k}^{\gamma}}{d E_{\gamma}}\left(E_{\gamma} \leftarrow E\right)=\sigma_{k}^{\gamma}(E) f_{k}\left(E_{\gamma} \leftarrow E\right),
$$

which obviously requires that

$$
\int_{0}^{E_{\gamma}^{\max }} f_{k}\left(E_{\gamma} \leftarrow E\right) d E_{\gamma}=1
$$

Any time a continuum representation is used for a given MT number in File 13, then the normalized energy distribution, $f_{k}\left(E_{\gamma} \leftarrow E\right)$, must be given in File 15 under the same MT number.

[^2]As a check quantity, the total photon production cross section,

$$
\sigma_{T O T}^{\gamma}(E)=\sum_{k=1}^{N K} \sigma_{k}^{\gamma}(E)
$$

is also tabulated for each MT number, unless only one subsection is present (i.e., $\mathrm{NK}=1$ ).

The structure of a section in File 13 is

$$
\begin{aligned}
& \text { ([MAT, 13, MT/ZA, AWR; b, b; NK, b]HEAD } \\
& \text { [MAT, 13, MT/b, } \left.b ; b ; b ; N R, N P / E_{\text {int }} / \sigma_{T O T}^{\gamma}(E)\right] T A B 1 * \\
& \left.\begin{array}{l}
\text { <subsection for } k=1 \\
\langle\text { subsection for } k=2
\end{array}\right\rangle \\
& \text { • } \\
& \text { - } \\
& \text { - } \\
& \text { <subsection for } k=N K \text { > } \\
& \text { [MAT, 13, 0/b, b; b, b; b, b]SEND }
\end{aligned}
$$

and the structure of each subsection is
[MAT, 13, MT/EG $\left.{ }_{k}, E S_{k} ; b, L F ; N R, N P / E_{\text {int }} / \sigma_{k}^{\gamma}(E)\right] T A B 1$,
where
$E S_{k} \equiv$ the energy of the level from which the photon originates. If the level is unknown or if a continuous photon spectrum is produced, then $E S_{k}=0.0$ should be used.
$E G_{k} \equiv$ the photon energy. For a continuous photon energy distribution, $E G_{k}=0.0$ should be used.
LF $\equiv$ the photon energy distribution law number, which presently has only two values defined: LF $=1$, a normalized tabulated function (in File 15), and $=2$, a discrete photon energy.

[^3]
### 2.2.2. File 13 Procedures

1. The subsections are given in decreasing magnitude of $E G_{k}$.
2. The convention is that the subsection for the continuum photons, if present, is last. In this case, $E G_{N K} \equiv 0.0$.
3. The values of $E G_{k}$ should be consistent to within four significant figures with the corresponding $E G_{k}$ values in File 14.
4. $E S_{k}$ is the energy of the level from which the photon originates, if known. Otherwise $E S_{k} \equiv 0.0$.
5. If capture and fission resonance parameters are given in File 2, the corresponding photon production should be given by use of Option 1 of File 12, instead of as photon production cross sections.
6. The total photon production cross section table, $\sigma_{T O T}^{\gamma}(E)$, should exactly span the same energy range as the combined energy range of all the $\sigma_{k}^{Y}(E)$. Within that range,

$$
\sigma_{\mathrm{TOT}}^{\gamma}(E)=\sum_{k=1}^{\mathrm{NK}} \sigma_{k}^{\gamma}(E)
$$

should hold within four significant figures. If only one energy distribution is given, either discrete or continuous ( $\mathrm{NK}=1$ ), then the TABl record for the $\sigma_{\text {TOT }}^{\gamma}(E)$ table is simply deleted.
7. The join of all sections in Files 12 and 13 combined should represent the photon production data with no redundancy. For example, MT $=4$ cannot include any photons given elsewhere under MT $=51$ through 91.
8. The limit on the number of energy points in a tabulation for any photon production subsection is 1000 . This is an upper limit; in practice, the minimum number of points possible should be used. If there is extensive structure, the use of File 12 should be seriously considered, because yields are normally much smoother functions of incident neutron energy than cross sections.
9. The limit on the number of interpolation regions is 10 .
10. Tabulations of nonthreshold data should normally cover at least the energy range $10^{-5} \mathrm{eV} \leq \mathrm{E} \leq 2 \times 10^{7} \mathrm{eV}$, where practical. Threshold data should be given from threshold energy up to $2 \times 10^{7} \mathrm{eV}$, where practical.

### 2.3. File 14: Photon Angular Distributions

The purpose of File 14 is to provide a means for representing the angular distributions of secondary photons produced in neutron interactions. Angular distributions should be given for each discrete photon and photon continuum that appears in Files 12 and 13, even if the distributions are isotropic.

The structure of File 14 is, with the exception of isotropic flags, closely analogous to that of File 4. Angular distributions for a specific reaction type (MT number) are given for a series of incident neutron energies in order of increasing neutron energy. The energy range covered should be the same as that for the data given under the corresponding reaction type in File 12 or File 13. The data are given in ascending order of MT number.

The angular distributions are expressed as normalized probability distributions, that is,

$$
\int_{-1}^{1} p_{k}(\mu, E) d \mu=1
$$

where $p_{k}(\mu, E)$ is the probability that an incident neutron of energy $E$ will result in a particular discrete photon or photon energy continuum (specified by $k$ and MT number) being emitted into unit cosine about an angle whose cosine is $\mu$. Because the photon angular distribution is assumed to have azimuthal symmetry, the distribution may be represented as a Legendre series expansion,

$$
\begin{aligned}
p_{k}(\mu, E) & =\frac{2 \pi}{\sigma_{k}^{\gamma}(E)} \frac{d \sigma_{k}^{\gamma}}{d \Omega}(\Omega, E) \\
& =\sum_{\ell=0}^{N L} \frac{2 \ell+1}{2} a_{\ell}^{k}(E) P_{\ell}(\mu),
\end{aligned}
$$

where
$\mu \equiv$ cosine of the reaction angle.
$E \equiv$ energy of the incident neutron in the laboratory system.
$\sigma_{k}^{Y}(E) \equiv$ photon production cross section for the discrete photon or photon continuum specified by $k$, as given in either File 13 or the combination of Files 2, 3, and 12.

$$
\begin{aligned}
\ell & \equiv \text { order of the Legendre polynomial. } \\
\frac{d \sigma_{k}^{\gamma}}{d \Omega} & \equiv \text { differential photon production cross section in barns/steradian. } \\
a_{\ell}^{\mathrm{k}}(\mathrm{E}) & \equiv \text { the } \ell \text { th Legendre coefficient associated with the discrete photon or } \\
& \text { photon continuum specified by } \left.k \text {. (It is understood that } a_{0}^{k}(E) \equiv 1.0 .\right)
\end{aligned}
$$

$$
a_{\ell}^{k}(E)=\int_{-1}^{1} p_{k}(\mu, E) P_{\ell}(\mu) d \mu
$$

Angular distributions may be given in File 14 by tabulating as a function of incident neutron energy either the normalized probability distribution function, $p_{k}(\mu, E)$, or the Legendre polynomial expansion coefficients, $a_{l}^{k}(E)$. Provision is made in the format for simple flags to denote isotropic angular distributions, either for a block of individual photons within a reaction type or for all photons within a reaction type taken as a group.

It is important to note that File 14 assumes separability of the photon energy and angular distributions for the continuous spectrum. If this is not the case, File 16 (analogous to File 6) must be used instead of Files 14 and 15. (Note that File 14 implicitly specifies an energy-angle distribution for discrete photons, and, therefore, File 16 is required only for the continuous spectrum.)

### 2.3.1. File 14 Format

As usual, sections are ordered by increasing reaction-type (MT) numbers.
The following definitions are required.
$\operatorname{LTT}=1$, data are given as Legendre coefficients, where $a_{0}^{k}(E) \equiv 1.0$ is understood.
$=2$, data are given as a tabulation.
LI $=0$, distribution is not isotropic for all photons from this reaction type (but may be for some photons),
$=1$, distribution is isotropic for all photons from this reaction type.
NE $\equiv$ number of neutron energy points given in a TAB2 record.
NI 三 number of isotropic photon angular distributions given in a section (MT number) for which $L I=0$, i.e., a section with at least one anisotropic distribution.
$\mathrm{NL}_{\mathrm{j}} \equiv$ highest value of $\ell$ required at each neutron energy $\mathrm{E}_{\mathrm{i}}$.
a. $\mathrm{LI}=1$ : Isotropic Distribution

If $\mathrm{LI}=1$, then all photons for the reaction type (MT) in question are assumed to be isotropic. This is a flag that the processing code can sense, thus avoiding the necessity to enter needless isotropic distribution data in the file. In this case, the section is composed of a HEAD card and a SEND card, as follows:

$$
\left\{\begin{array}{l}
{[\operatorname{MAT}, 14, \mathrm{MT} / \mathrm{ZA}, \mathrm{AWR} ; \mathrm{LI}=1, \mathrm{~b} ; \mathrm{NK}, \mathrm{~b}] \text { HEAD }} \\
{[\mathrm{MAT}, 14,0 / \mathrm{b}, \mathrm{~b} ; \quad \mathrm{b}, \mathrm{~b} ; \mathrm{b}, \mathrm{~b}] \text { SEND } .}
\end{array}\right.
$$

b. LI = 0: Anisotropic Distribution

If $\mathrm{LI}=0$, there are two possible structures for a section, depending upon the value of LTT, but the section always starts with a HEAD record of the form:
[MAT, 14, MT/ZA, AKR; LI=0, LTT; NK, NI]HEAD .

1. LTT $=1$ : Legendre Coefficient Representation


The structure of each record in the first block of NI subsections, which is for the NI isotropic photons, is
[MAT, 14, MT/EG,$\left.E_{k} ; b, b ; b, b\right] C O N T$ •
That is, there is just one CONT record for each isotropic photon. (The set of CONT records is empty if NI = O.) The subsections are ordered in decreasing magnitude of $E G_{k}$, and the continuum, if present and isotropic, appears last with $E G_{k} \equiv 0.0$.

This block of NI subsections is then followed by a block of NK-NI subsections for the anisotropic photons in decreasing magnitude of $E G_{k}$ (photon energy). The continuum, if present and anisotropic, appears last with $E G_{k} \equiv 0.0$. The structure for the last NK-NI subsections is

$$
\begin{aligned}
& {\left[\text { MAT, 14, MT/EG }{ }_{k}, E S_{k} ; b, b ; N R, N E / E_{\text {int }}\right] T A B 2} \\
& \text { [MAT, 14, MT/ b, } \left.E_{1} ; b, b ; L_{1}, b / a \frac{k}{\ell}\left(E_{1}\right)\right] \text { LIST } \\
& \text { [MAT, 14, MT/ b, } \left.\mathrm{E}_{2} ; \mathrm{b}, \mathrm{~b} ; \mathrm{NL}_{2}, \mathrm{~b} / \mathrm{a}_{\ell}^{\mathrm{k}}\left(\mathrm{E}_{2}\right)\right] \text { LIST } \\
& \text { [MAT, 14, MT/ } \left.b, E_{N E} ; b, b ; \operatorname{NL}_{N E}, b / a_{\ell}^{k}\left(E_{N E}\right)\right] L I S T \text {. }
\end{aligned}
$$

Note that lists of the $a_{\ell}^{k}(E)$ start at $\ell=1$ because $a_{0}^{k}(E) \equiv 1.0$ is always understood.
ii. $\operatorname{LTT}=2:$ Tabulated Angular Distributions

The structure of a section for $\mathrm{LI}=0$ and $\mathrm{LTT}=2$ is


The structure of the first block of NI subsections (where NI may be zero) is the same as for the case of a Legendre representation, i.e., it consists of one CONT record for each of the NI isotropic photons
in decreasing magnitude of $E G_{k}$ (photon energy). The continuum, if present and isotropic, appears last with $E G_{k} \equiv 0.0$.
The structure of the first NI subsections is
[MAT, 14, MT/EG $\left.{ }_{k}, E S_{k} ; b, b ; b, b\right] C O N T$.
This block of NI subsections is then followed by a block of NK-NI subsections for the anisotropic photons, again in decreasing magnitude of $E G_{k}$ with the continum, if present and anisotropic, appearing last with $E G_{k} \equiv 0.0$. The structure of the last NK-NI subsections is
[MAT, 14, MT/EG $\left.{ }_{k}, E S_{k} ; b, b ; N R, N E / E E_{\text {int }}\right]$ TAB2
[MAT, 14, MT/ b, $\mathrm{E}_{1} ; \mathrm{b}, \mathrm{b}$; NR, NP/ $\left.\mu_{\text {int }} / \mathrm{p}_{\mathrm{k}}\left(\mathrm{H}, \mathrm{E}_{1}\right)\right]$ TAB1
[MAT, 14, MT/ b, $\left.E_{2} ; b, b ; N R, N P / \mu_{\text {int }} / p_{k}\left(\mu, E_{2}\right)\right] T A B 1$
[MAT, 14, MT/ b, $E_{N E} ; \mathrm{b}, \mathrm{b}$; NR, NP/ $\left.\mu_{\text {int }} / \mathrm{p}_{\mathrm{k}}\left(\mu, \mathrm{E}_{\mathrm{NE}}\right)\right] \mathrm{TAB1}$.

### 2.3.2. File 14 Procedures

1. The subsections are given in decreasing magnitude of $E G_{k}$ within each of the isotropic and anisotropic blocks.
2. The convention is that the subsection for the continuous photon spectrum, if present, appears last in its block. In this case, $E G_{N K} \equiv 0.0$.
3. The values of $E G k$ should be consistent within four significant figures with the corresponding $E G_{k}$ values in File 12 or 13 . In the case of File 12 , Option 2 (transition probability arrays), the values of $E G_{k}$ are implicitly determined by the level energies.
4. $E S_{k}$ is the energy of the level from which the photon originates, if known. Otherwise, $E S_{k} \equiv 0.0$ (as is always the case for the continuum).
5. Data should not appear in File 14 for photons that do not have production data given in File 12 or 13. Conversely, every photon appearing in File 12 or 13 must have an angular distribution given in File 14. The neutron energy range over which the angular distributions are given should be the same as that for which the photon production data are given in File 12 or 13.
6. For $L T T=1$ (Legendre coefficients), the value of NL should be the minimum number of coefficients that will reproduce the angular distribution with sufficient accuracy and be positive everywhere. In all cases, NL should be an even number less than or equal to 20 .
7. The $T A B 1$ records for the $p_{k}\left(\mu, E_{i}\right)$ within a subsection are given in increasing order of neutron energy, $E_{i}$.
8. The tabulated probability functions, $p_{k}\left(\mu, E_{i}\right)$, should be normalized within four significant figures (to unity).
9. The interpolation scheme for $p_{k}(\mu, E)$ with respect to $E$ must be linearlinear or log-linear (INT $=2$ or 3 ) to preserve normality of the interpolated distributions. It is recommended that the interpolation in $\mu$ be linear-linear (INT = 2) .
10. For $L I=1$ (isotropic distribution), the parameter NK is a count of the number of photons in that section and should be given consistent with the NK values in Files 12 and 13. This parameter could also be determined independently from Files 12 and 13, but it is useful in File 14 for the "physics" checking code. 6
11. The minimum amount of data should be used that will accurately represent the angular distribution as a function of both $\mu$ and $E$.
12. If all photons for a reaction type (MT number) are isotropic, the LI $=1$ flag should be used. The use of LI $=0$ and NI $=N K$ is strongly discouraged. Likewise, isotropic distributions should not be entered explicitly as a tabulation or as a Legendre expansion with $a_{\ell}^{k}(E) \equiv 0, \ell \geq 1$.

### 2.4. File 15: Continuous Photon Energy Spectra

File 15 provides a means for representing continuous energy distributions of secondary photons. The energy distributions are expressed as normalized probability distributions. The energy distribution of each photon continuum that occurs in Files 12 and 13 should be specified in File 15 over the same neutron energy range as used in Files 12 and 13 . File 15 is divided into sections with each section giving the data for a particular reaction type (MT number), and the sections are ordered by increasing MT number.

The energy distributions, $f\left(E_{\gamma} \leftarrow E\right)$, have the units of $\mathrm{eV}^{-1}$ and are normalized such that

$$
\int_{0}^{E_{Y}^{\max }} f\left(E_{Y} \leftarrow E\right) d E_{Y}=1
$$

where $E_{\gamma}^{\max }$ is the maximum possible secondary photon energy and its value depends on the incoming neutron energy as well as the particular nuclei involved.* The energy distributions $f\left(E_{\gamma} \leftarrow E\right)$ can be broken down into the weighted sum of several different normalized distributions in the following manner

$$
f\left(E_{\gamma} \leftarrow E\right)=\sum_{j=1}^{N C} p_{j}(E) q_{j}\left(E_{\gamma} \nleftarrow E\right) \quad(e V)^{-1}
$$

where
$N C \equiv$ the number of partial distributions used to represent $f\left(E_{\gamma} \leftarrow E\right)$. $q_{j}\left(E_{\gamma} \leftarrow E\right) \equiv$ the jth normalized partial distribution in the units $\mathrm{eV}^{-1}$. $\begin{aligned} p_{j}(E) \equiv & \text { the probability or weight given to the } j \text { th partial distribution, } \\ & q_{j}\left(E_{\gamma}+E\right) .\end{aligned}$

[^4]The following normalization condition is imposed.

$$
\int_{0}^{E_{\gamma}^{\max }} q_{j}\left(E_{\gamma}+E\right) d E_{\gamma}=1
$$

Thus,

$$
\sum_{j=1}^{N C} p_{j}(E)=1
$$

The absolute energy distribution cross section, $\sigma^{\gamma}\left(E_{\gamma} \leftarrow E\right)$, can then be constructed from the expression

$$
\sigma^{\gamma}\left(E_{\gamma}+E\right)=\sigma^{\gamma}(E) f\left(E_{\gamma}+E\right) \quad(b / e V)
$$

where $\sigma^{\gamma}(E)$ is the integrated cross section for the continuum given either directly in File 13 or through the combination of Files 2, 3, and 12.

The system used to represent continuous photon energy distributions in File 15 is similar to that used in File 5. At present, however, there is only one continuous distribution law activated for File 15, i.e.,

$$
q_{j}\left(E_{\gamma} \leftarrow E\right)=g\left(E_{\gamma} \leftarrow E\right),
$$

where $g\left(E_{\gamma} \leftarrow E\right)$ represents an arbitrary tabulated function. In the future, new laws (for example, the fission ganma-ray spectrum) may be added.

### 2.4.1. File 15 Format

The structure of a section is

$$
\left\{\begin{array}{c}
{[\text { MAT, } 15, \text { MT /RA, AWR; } b, b ; N C, b] H E A D} \\
\langle\text { subsection for } j=1\rangle \\
\langle\text { subsection for } j=2\rangle \\
\cdot \\
\cdot \\
\langle\text { subsection for } j=N C\rangle
\end{array}\right.
$$

For $L F=1$, the structure of a subsection is

At present, there is no other continuous energy distribution law defined. Therefore, formats for other laws remain to be defined in the future, but their structures will probably closely parallel those in File 5 for $L F=5$, 7, 9, and 10.

### 2.4.2. File 15 Procedures

1. Photon energies, $E_{\gamma}$, within a subsection are given in increasing magnitude.
2. The TABI records for the $g\left(E_{\gamma} \leftarrow E_{i}\right)$ within a subsection are given in increasing order of nuetron energy, $E_{i}$.
3. The tabluated functions, $g\left(E_{\gamma} \leftarrow E_{i}\right)$, should be normalized to unity within four significant figures.
4. The interpolation scheme for $p_{j}(E)$ must be either linear-linear or loglinear (INT $=1,2$, or 3) to preserve probabilities upon interpolation. Likewise, the interpolation scheme for $g\left(E_{\gamma} \leftarrow E\right)$ must be linear-linear or log-linear with respect to $E$.
5. The neutron energy mesh should be a subset of that used for the $y_{N K}(E)$ tabulation in File 12 or for the $\sigma_{\mathrm{NK}}^{\gamma}(\mathrm{E})$ tabulation in File 13, and the energy ranges must be identical. However, the neutron energy mesh for the $p_{j}(E)$ need not be the same as for the $g\left(E_{\gamma}+E\right)$, as long as they span the same range.
6. For an MT number which appears in both File 12 and File 13, a continuous photon energy distribution ( $L F=1$ ) can appear in only one of those files. Otherwise the distribution as given in File 15 could not in general be uniquely associated with a corresponding multiplicity or production cross section.
7. The minimum amount of data should be used that will accurately represent the energy distribution as a function of both $E_{\gamma}$ and $E$. However, caution must be exercised not to use too coarse a mesh for $E$, even if the distributions are slowly varying functions of $E$. This is due to the fact that the interpolated distribution will always have a nonzero component up to the maximum energy to which either of the original distributions has a nonzero component.
8. The limit on the number of neutron energy points for either $p_{j}(E)$ or $g\left(E_{\gamma} \leftarrow E\right)$ is 200. The limit on the number of photon energy points for $g\left(E_{Y}+E\right)$ is 1000 .

### 2.5. File 16: Photon Energy-Angle Distributions

If the photon energy dependence and angular dependence of the cross section for production of the continuous spectrum are not separable, then a file analogous to File 6 is required instead of Files 14 and 15. The energy-angle distributions for discrete photons are completely determined in File 14 and should never appear in File 16. A knowledge of the formats for Files 14 and 15 is assumed in this discussion.

Consider a nonnormalized energy-angle distribution function $F\left(E_{\gamma} \leftarrow E, \mu\right)$, where the angular dependence is normalized such that

$$
\int_{-1}^{1} F\left(E_{\gamma}+E, \mu\right) d \mu=y\left(E_{\gamma}+E\right)
$$

Then the multipiicity (yield) can be separated out, leaving a function, $h\left(E_{\gamma}+E, \mu\right)$, normalized in both $E_{\gamma}$ and $\mu$.

$$
F\left(E_{\gamma} \leftarrow E, \mu\right)=y(E) h\left(E_{\gamma} \leftarrow E, \mu\right)
$$

The differential photon production cross section is then obtained from

$$
\frac{\partial^{2} \sigma\left(E_{\gamma} \leftarrow E, \mu\right)}{\partial E_{\gamma} \partial \mu}=\sigma(E) y(E) h\left(E_{\gamma} \leftarrow E, \mu\right) \quad \quad \text { (b-photons/eV) }
$$

where $\sigma(E)$ is the cross section for the reaction type being considered, as determined by Files 2 and 3.

As in File 14, the angular part of the distribution may be specified either in tabular form or as Legendre coefficients, $\eta_{\ell}\left(E_{\gamma}+E\right)$. The Legendre expansion is

$$
h\left(E_{\gamma} \leftarrow E, \mu\right)=\sum_{\ell=0}^{N L} \frac{2 \ell+1}{2} n_{\ell}\left(E_{\gamma} \leftarrow E\right) P_{\ell}(\mu)
$$

Now consider the structure of a section for the two possible forms of the angular distribution.

### 2.5.1. File 16 Format

Two options are allowed, corresponding to the options in File 14, i.e., the angular distribution can be represented by either Legendre coefficients (LTT = 1) or by tabulated angular distributions (LTT = 2).

LTT $=1$ : Legendre Coefficient Representation. In this option, the Legendre coefficients are tabulated as functions of both incident neutron energy and photon energy. The structure of a section for LTT $=1$ is

```
([MAT, 16, MT/ZA, AWR; b, LTT=1; b, b]HEAD
    [MAT, 16, MT/ b, b; b, b; NL, b]CONT
            \(\langle\) subsection for \(\ell=0\rangle\)
            \(\langle\) subsection for \(\ell=1\rangle\)
            < subsection for \(\ell=N L\) >
! [MAT, 16, 0/b, b; b, b; b, b]SEND.
```

The subsections contain the energy distributions, and the structure of a subsection is identical to the structure of a section for a continuous energy distribution (File 15), with the following exceptions
a. The SEND record is deleted.
b. The HEAD record is changed to read
[MAT, $16, \mathrm{MT} / \mathrm{b}, \mathrm{b} ; \mathrm{b}, \mathrm{b}$; NC, b]CONT •
c. $g\left(E_{\gamma} \leftarrow E\right)$ is replaced by $\eta_{\ell}\left(E_{Y} \leftarrow E\right)$.

LTT $=2:$ Tabulated Angular Distribution. In this option, the subsections consist of tabulations for $h\left(E_{\gamma} \leftarrow E, \mu_{m}\right), m=1(1) N A$. The structure of a section for LTT $=2$ is

```
[ \([\mathrm{MAT}, 16, \mathrm{MT} / \mathrm{ZA}, \mathrm{AWR} ; \mathrm{b}, \mathrm{LTT}=2\); b, b]HEAD
[MAT, 16, MT/b, b; b, b; NR, NA/ \(\mu_{\text {int }}\) ]TAB2
        <subsection for \(m=1\rangle\)
        <subsection for \(m=2\) 〉
                        .
        < subsection for \(m=N A>\)
    [MAT, 16, 0/b, b; b, b; b, b]SEND •
```

As in the case of the Legendre coefficient representation, this section for a tabulation contains subsections. The structure of a subsection is identical to the structure of a section for File 15, with the following exceptions
a. The SEND record is deleted.
b. The HEAD record is changed to read
[MAT, 16, MT/b, $\left.\mu_{m} ; b, b ; N C, b\right] C O N T$.
c. $g\left(E_{\gamma} \leftarrow E\right)$ is replaced by $g\left(E_{\gamma} \leftarrow E, \mu_{m}\right)$, where each subsection is for a particular value of $\mu_{m}, m=1(1) N A$.

### 2.5.2. File 16 Procedures

The procedures for this file are the same as those listed for Files 14 and 15, where applicable.

## 3. PHOTON INTERACTION

Photon interaction data are divided into five files, the first four of which are analogous to Files 3 through 6.

| $\frac{\text { File }}{23}$ | Description |
| :--- | :--- |
| 24 | "Smooth" cross sections |
| 25 | Secondary angular distributions |
| 26 | Secondary energy distributions |
| 27 | Secondary energy-angle distributions |
| Form factors for coherent and incoherent scattering. |  |

As with the photon production data files, the photon interaction data formats were kept as closely parallel as possible to those for the neutron data files of the same number (modulo 10). This facilitates the use of existing retrieval routines in processing codes for photon interaction data (as in CHECKER). ${ }^{4}$ Also, the format is kept as consistent as possible with the United Kingdom Atomic Energy Authority data file. ${ }^{5}$ For Compton scattering at higher energies ( $\gtrsim 1 \mathrm{MeV}$ ), the energy and angular distribution files would not normally be used because a simple analytical representation of these distributions is available. Also, provision is made for the entry of coherent scattering form factors as well as incoherent scattering form factors (incoherent scattering functions). The secondary energy and angular distribution files can be used for both photon secondaries or particulate secondaries (e.g., photoneutrons).

Procedures are given for Files 23 and 27, but none will be given for Files 24,25 , and 26 until those files are activated. There are, at present, no data in these files.
3.1. File 23: "Smooth" Cross Sections

This file is for the integrated photon interaction cross sections, including those usually called microscopic attenuation or energy-deposition coefficients, as well as photonuclear reaction cross sections. The reaction type (MT) numbers for photon interaction are in the 500 and 600 series. Several common photon interactions have been assigned MT numbers as follows

MT Reaction Description
501 Total
502 Coherent scattering
504 Incoherent scattering
515 Pair production, electron field
516 Pair production, nuclear and electron field
(i.e., pair plus triplet production)

Pair production, nuclear field
Photofission ( $\gamma, F$ )
Photoneutron ( $\gamma, \mathrm{n}$ )
Total photonuclear Photoelectric

Photon cross sections, such as the total cross section, coherent elasticscattering cross section, and incoherent (Compton) cross section, are given in File 23, which is essentially the same in structure as File 3. These data are given as a function of energy, $E_{\gamma}$, where $E_{\gamma}$ is the energy of the incident photon (in eV). The data are given as energy-cross-section pairs. An interpolation scheme is given that specifies the energy variation of the cross section for photon energies between a given energy point and the next higher energy point. The photon cross sections are given in one or more energy ranges. Within any one energy range, the interpolation scheme is unchanged. The interpolation scheme may change from one to another energy range.

File 23 is divided into sections with each section containing the data for a particular reaction type (MT number). The sections are ordered by increasing reaction type numbers.

### 3.1.1. File 23 Format

The format is almost identical to that of File 3, as follows
[MAT, 23, MT/ZA, AWR; b, b; b, b]HEAD
[MAT, 23, MT/b, $b ; b, b ; N R, N P / E_{\gamma}$ int $\left./ \sigma\left(E_{\gamma}\right)\right] T A B 1$ [MAT, 23, $0 / \mathrm{b}, \mathrm{b} ; \mathrm{b}, \mathrm{b}$; b, b]SEND.

### 3.1.2. File 23 Procedures

1. Values are usually for elements, hence, except for monoisotopic elements, $Z A=Z \times 1000 ;$ also, AWR should be for the naturally occurring element.
2. Photoelectric edges will not be multivalued. That is, the edge will be defined by two energies different in the fourth or fifth significant figure.
3. The total pair production values are given for reaction type $\mathrm{MT}=516$. Reaction type 517 is reserved for the portion of the pair production cross section due to the nuclear field, i.e., excluding triplet production.
4. Interpolation is normally log-log (INT $=5$ ).
5. Kerma factor (energy deposition coefficients) libraries will normally be local because there is no universal definition. The application will determine whether annihilation or other radiation fractions are subtracted.
3.2. File 24: Secondary Angular Distribution

The structure of File 24 is identical to that for File 4 , so the pertinent discussion from Vol. I of this report will be reviewed here for convenience.

Secondary angular distributions are expressed as probability density functions, $p\left(\mu, E_{\gamma}\right)$. These functions can be represented either as a tabulation or as the Legendre coefficients, $f_{\ell}\left(E_{\gamma}\right)$, in

$$
\begin{aligned}
\frac{d \sigma\left(E_{\gamma}, \mu\right)}{d \mu} & =\sigma\left(E_{\gamma}\right) \sum_{\ell=0}^{N L} \frac{2 \ell+1}{2} f_{\ell}\left(E_{\gamma}\right) P_{\ell}(\mu) \\
f_{o}\left(E_{\gamma}\right) & =1.0
\end{aligned}
$$

Here, $\mu=\cos \theta$, where $\theta$ is the polar angle of scattering in either the center-of-mass (C) or the laboratory (L) system. The secondary may be either a photon (coherently scattered) or a particle (e.g., photoneutrons). For the case where the secondary distribution is for a photon, the laboratory system is always used.

### 3.2.1. File 24 Format

The format is identical to that for File 4 and will not be reproduced here (see Vol. I). However, for the case where the secondary distribution is for a photon, the LCT flag is not relevant, and the following arbitrary convention is adopted
a. $\operatorname{LCT}=1$ [data are given in the laboratory (L) system].
b. LVT $=0$ [transformation matrix is not given].
3.3. File 25: Secondary Energy Distributions

The structure of the analogous File 5 appears to be entirely adequate (see Vo1. I). Thus, the format will not be reproduced here, but will be adopted by reference to File 5.
3.4. File 26: Secondary Energy-Angle Distributions

The structure of the analogous File 6 appears to be entirely adequate (see Vol. I). Thus, the format will not be reproduced here but adopted by reference to File 6. The inclusion of File 26 (as well as Files 6 and 16) is, at the present stage of development of cross-section data, strictly pro forma.

### 3.5 File 27: Atomic Form Factors or Scattering Functions

The ENDF system for neutron and photon production data allows two alternatives for storing angular distribution data. One is by probability per unit $\cos \theta$ vs $\cos \theta$, and the other is by Legendre coefficients. Actually, neither of these is a "natural" method for photons. The natural method for storing photon distributions would be atomic form factors or incoherent scattering functions. These are discussed briefly below.
a. Incoherent Scattering. The cross section for incoherent scattering is given by

$$
\frac{\mathrm{d} \sigma_{i}}{\mathrm{~d} \mu}=Z K(q ; Z) \frac{\mathrm{d} \sigma_{c}}{\mathrm{~d} \mu}
$$

where $d \sigma_{C} / d \mu$ is the Klein-Nishina cross section, which can be written in closed form. The factor $K(q ; Z)$ is the incoherent scattering function, which is often symbolized as $S(q ; Z)$ in the literature. At high ( $\gtrsim 1 \mathrm{MeV}$ ) energies, $K$ approaches unity. The quantity $q$ is the momentum of the recoil electron (in units of $m_{0} c$,

$$
q=\alpha\left[1+\left(\frac{\alpha^{-}}{\alpha}\right)^{2}-2 \mu\left(\frac{\alpha^{-}}{\alpha}\right)\right]^{\frac{1}{2}}
$$

where

$$
\begin{aligned}
\alpha & =E_{\gamma} / m_{o} c^{2} \\
E_{\gamma}^{\prime} & =\text { scattered photon energy, and } \\
\mu & =\cos \theta .
\end{aligned}
$$

The angular distribution can then easily be calculated, given a table of $K(q ; Z)$. Because $K$ is a smoothly varying function of $q$, it can be represented by a reasonably small array of numbers. The quantities $K(q ; Z)$ will be tabulated as a function of $q$ in File 27. The user presumably will have subroutines available for calculating $q$ for energies and angles of interest and for calculating KleinNishina cross sections. He will then generate his cross sections for the appropriate cases by calculating $q^{\prime}$, looking up the appropriate values of $K$, and substituting them into the above formula.
b. Coherent Scattering. The coherent scattering cross section is given by

$$
\frac{d \sigma_{c o h}}{\mathrm{~d} \mu}=\pi r_{0}^{2} z^{2}\left(1+\mu^{2}\right) G(q ; z),
$$

where
$q=\alpha[2(1-\mu)]^{\frac{1}{2}}$, the recoil momentum of the atom (in units of $m_{o} c$ ), and $r_{o}=e^{2 / m_{o}} c^{2}$, the classical radius of the electron.
The quantity $G(q ; Z)$ is a form factor which is often symbolized as $F(q ; Z)$ in the literature. This quantity is also easily tabulated. At high ( $\gtrsim 1 \mathrm{MeV}$ ) energies, G approaches zero.

An alternative way of presenting the photon scattering data, then, would be to tabulate incoherent scattering functions and form factors. Users could then provide processing codes to generate the cross sections from this information. The calculation is quite straightforward and allows the user to generate all his scattering data from a relatively small table of numbers. The incoherent and coherent scattering data should always be presented as scattering functions and form factors, respectively, regardless of whether or not data are included in Files 24, 25, or 26.

### 3.5.1. File 27 Format

The structure of a section is then very similar to File 3 (and 23), as follows.
$\left\{\begin{array}{l}{[\text { MAT, 27, MT/ZA, AWR; b, b; b, b]HEAD }} \\ {\left[\text { MAT, 27, MT/b, } \mathrm{z} ; \mathrm{b}, \mathrm{b} ; \mathrm{NR}, \mathrm{NP} / \mathrm{q}_{\text {int }} / \mathrm{H}(\mathrm{q} ; \mathrm{z})\right] \text { TAB1 }} \\ {[\text { MAT, 27, } 0 / \mathrm{b}, \mathrm{b} ; \mathrm{b}, \mathrm{b} ; \mathrm{b}, \mathrm{b}] \text { SEND. }}\end{array}\right.$
The general symbol $H(q ; Z)$ is used for either $G(q ; Z)$ or $K(q ; Z)$ for coherent and incoherent scattering, respectively.

### 3.5.2. File 27 Procedures

1. Values of $G(q ; Z)$ and $K(q ; Z)$ should be entered for the respective entire energy range for which integrated coherent and incoherent cross sections are given in File 23. This is true even though the value may be 0.0 or 1.0 , respectively, over most of the (higher) energy range.
2. The value of $Z$ is entered in floating-point format.

## ACKNOWLEDGMENT

The editor is grateful to the many people who contributed suggestions or written proposals concerning suitable formats and procedures for photon data. Also, appreciation is due those who have reviewed the report in its various draft forms and whose comments have helped shape this final version. Because of the many iterations to which this report has been subjected, it is difficult, in most cases, to identify specific sections with individual contributors. However, the listed contributors have, in some way, directly influenced the structure of the format itself. It is hoped that those who see their suggestions or comments reflected in this report will not feel slighted by lack of specific acknowledgment. The thorough and expert reviews by H. Clyde Claiborne, Robert W. Roussin, and Robert E. Seamon certainly deserve special recognition.

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## APPENDIX A

GLOSSARY OF VARIABLE NAMES AND SYMBOLS

| $A_{j, i}$ | $\left(T P_{j, i}\right)\left(G P_{j, i}\right)$, the probability of photon emission by a direct transition ${ }^{\prime}$ rom level $j$ to level $i$. |
| :---: | :---: |
| $\mathrm{a}_{\ell}^{\mathrm{k}}(\mathrm{E})$ | The (neutron) energy-dependent Legendre coefficients in the Legendre expansion of the $k$ th photon angular distribution. |
| AWR | The atomic weight ratio of the target nucleus, i.e., the ratio of the atomic mass to the mass of the neutron. |
| b | A blank field or the unit "barns." |
| E | Incident neutron energy (eV). |
| EY | Photon energy (eV). |
| $E G_{k}$ | The photon energy for the kth subsection within a reaction type (eV). |
| $E S_{k}, E S_{i}$ | The energy of the level from which a photon originates. In a transition probability array, $\mathrm{ES}_{i}$ is the energy of the ith level (eV). |
| $F\left(E_{\gamma} \leftarrow E, \mu\right)$ | An energy-angle distribution function for photon production (photons/eV). |
| $\mathrm{f}_{\mathrm{k}}\left(\mathrm{E}_{\gamma} \leftarrow \mathrm{E}\right)$ | A normalized (to unity) photon energy distribution (or probability density) function at incident neutron energy $E$ for the kth subsection within a reaction type ( $\mathrm{eV}^{-1}$ ). |
| $\mathrm{f}_{\ell}\left(E_{\gamma}\right)$ | The (photon) energy-dependent Legendre coefficients in the Legendre expansion of the secondary angular distribution. |
| $g\left(E_{\gamma} \leftarrow E\right)$ | A particular class of the functions $q_{j}\left(E_{\gamma} \leftarrow E\right)$ in File 15; those which are tabulated $\left(\mathrm{eV}^{-1}\right)$. |
| $G(q ; Z)$ | The form factor for coherent photon scattering. |
| $\mathrm{GP}_{j, i} \equiv \mathrm{GP}_{i}$ | The conditional probability of photon emission in a direct transition from level $j$ to level $i, i<j$. |
| $h\left(E_{\gamma} \leftarrow E, \mu\right)$ | A normalized (to unity) energy-angle distribution function for photon production ( $\mathrm{eV}^{-1}$ ). |
| H(q; Z ) | A general symbology for a form factor or incoherent scattering function; either $G(q ; Z)$ or $K(q ; Z)$, respectively. |


| $x_{\text {int }}$ | The interpolation table for variable x (cf. Ref. 2, pp. 3.8, 4.5 and 4.6). |
| :---: | :---: |
| $\mathrm{K}(\mathrm{q} ; \mathrm{z})$ | The incoherent scattering function for incoherent photon scattering. |
| LF | The photon energy distribution law number. |
| LG | The transition probability array flag for distinguishing between doublet and triplet arrays in File 12. |
| LI | The isotropy flag in File 14. |
| L $\varnothing$ | The option flag to determine whether multiplicities or transition probability arrays are to be given in File 12. |
| LCT | The option flag to determine whether secondary angular or energy data are given in the laboratory or center-of-mass system. |
| LFS | The final state number (Ref. 2, p. 7.1). |
| LTT | The option flag to determine whether angular distributions are to be given as Legendre coefficients or as tabulations. |
| LVT | The flag to determine whether or not a transformation matrix is given for converting between the laboratory and the center-ofmass systems. |
| $\mathrm{m}_{0}$ | The mass of the electron (ev-sec ${ }^{2} / \mathrm{cm}^{2}$ ). |
| MF | The file number (positive integer < 100). |
| MT | The reaction type number (positive integer < 1000). |
| MAT | The material number (positive integer $<10,000$ ). |
| NA | A count of the number of subsections in File 16 for an energyangle distribution that has tabulated angular distributions, (for $\left.h\left(E_{\gamma} \leftarrow E, \mu_{m}\right), m=1(1) N A\right)$. |
| NC | A count of the number of subsections in File 15 for normalized continuous photon energy distributions, (for $q_{j}\left(E_{\gamma} \leftarrow E\right)$, $j=1(1) N C)$. |
| NE | A count of the number of energy tabulation points in a TAB2 record. |
| NI | A count of the number of isotropic photon angular distributions given in a section for which LI = 0; i.e., a section with at least one anisotropic distribution. |


| NK | A count of the number of subsections in Files 12 and 13 for photon production multiplicities or cross sections, i.e., for $y_{k}(E)$ or $\sigma_{k}^{Y}(E)$, respectively, $k=1(1) N K$. |
| :---: | :---: |
| NL | A count of the order of the Legendre expansion in Files 14, 16, 24 , and 26 , i.e., $\ell=0(1) \mathrm{NL}$. |
| NP | A count of the number of points in a tabulation in a TAB1 record. |
| NR | A count of the number of interpolation regions in a tabulation in either a TABl or a TAB2 record. |
| NS | A count of the number of levels (including the ground state) below a level for which a transition probability array is being given in File 12. Thus, NS is the order number of the level being considered. |
| NT | A count of the number of transitions for which data are given in a LIST record in File 12 , Option $2(L \emptyset=2)$. |
| $\mathrm{p}_{\mathrm{k}}(\mu, E)$ | A normalized (to unity) photon angular distribution function at incident neutron energy $E$ for the $k$ th subsection. |
| $p_{j}(E)$ | The occurrence probability of the $j$ th piecewise continuous and normalized photon energy distribution function in File 15, $q_{j}\left(E_{Y} \leftarrow E\right)$ |
| q | The momentum of the recoil electron after an incoherent photon scattering event or the momentum of the recoil atom after a coherent photon scattering event ( $m_{0} c^{2}$ units). |
| Q | The reaction energy balance; the usual "Q-value" (eV). |
| $\mathrm{q}_{j}\left(\mathrm{E}_{Y} \leftarrow E\right)$ | A normalized (to unity) piecewise continuous photon energy distribution function that comprises the $j$ th component of the total such function $\left(\mathrm{eV}^{-1}\right)$. |
| $T P_{j, i}$ | $\equiv \mathrm{TP}_{1}$, the probability of a direct transition from level j to level 1 , $i<j$. |
| $y\left(E_{\gamma} \leftarrow E\right)$ | A total photon yield (multiplicity) energy distribution function for photon production (photons/eV). |
| Y (E) | The total photon yield (multiplicity) for a reaction type (photons). |
| $y_{k}(E)$ | The partial photon yield (multiplicity) for the kth subsection within a reaction type (photons). |
| Z | The atomic number of a material in File 27. |
| ZA | $=1000.0 \times Z+A$, where $A$ is the atomic mass number. For polyisotopic materials, $A$ is set to 0.0 . |


| $\alpha$ | $=\frac{E_{\gamma}}{m_{0} c^{2}} \text {, a dimensionless energy unit. }$ |
| :---: | :---: |
| $\eta_{\chi}\left(E_{Y} \leftarrow E\right)$ | The neutron- and photon-energy-dependent Legendre coefficients in the Legendre expansion of the photon energy-angle distribution ( $\mathrm{eV}^{-1}$ ). |
| $\mu$ | $=\cos \theta$, the cosine of the polar angle of scattering. |
| $\sigma_{c}$ | The integrated Klein-Nishina cross section (b). |
| $\sigma_{\mathrm{coh}}$ | The integrated coherent photon scattering cross section (b). |
| $\sigma_{i}$ | The integrated incoherent photon scattering cross section (b). |
| $\sigma_{\mathrm{TOT}}^{\gamma}(E)$ | The total photon production cross section in File 13 for a reaction type MT (b). |
| $\sigma_{k}^{Y}(E)$ | The partial photon production cross section in File 13 for the kth photon energy distribution subsection (b). |
| $\sigma_{j, i}{ }^{(E)}$ | The cross section to produce the photon that arises from decay of level $j$ to level $i, i<j$ (b). |
| $\sigma_{k}(E)$ | The neutron interaction cross section for inelastic scattering to level $k$ (b). |

APPENDIX B

BCD CARD-IMAGE FORMATS

## FILE 12: Multiplicities and Transition Probability Arrays

$$
\mathrm{L} \phi=1 \text { (Option } 1 \text {, Multiplicities) }
$$

NK the total number of subsections for a reaction type section (MT number), i.e., the number of discrete photons, plus one if a photon continuum is given.

NR the number of interpolation ranges given.
NP the total number of neutron energy (E) points in the tabulation to follow, i.e., the number of energy-multiplicity pairs.
$\mathrm{NBT}_{i}, \mathrm{INT}_{i}$ the interpolation scheme for the ith interpolation range.
$E_{i}$ the neutron energy for the ith pair in the tabulation.
$\underline{Y\left(E_{i}\right)}$ the total yield for all subsections (not given if $N K=1$ ).

$$
Y(E)=\sum_{k=1}^{N K} y_{k}(E)
$$

$y_{k}\left(E_{i}\right)$ the partial multiplicity or yield for the $k$ th subsection within a reaction type section, i.e., for a particular discrete photon or a photon continuum. In the case of a continuum, it is the integrated (over photon energy) multiplicity; the normalized photon energy distribution is given in File 15.
$E S_{k}$ the energy of the level from which the photon originates. If the - level is unknown or if a continuous photon spectrum is produced, then $\mathrm{ES}_{\mathrm{k}} \equiv 0.0$ should be used.
${ }^{E G}{ }_{k}$ the photon energy. For a continuous photon energy distribution, $E G_{k} \equiv 0.0$ should be used.

LF the photon energy distribution law number, which presently has only two values defined:
LF $=1$, a normalized tabulated function (in File 15), and $=2$, a discrete photon energy.
bcd card-image format for file 12
uutiplicities and transition probability array

| $\begin{gathered} \text { Field } \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 2 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ \hline \end{gathered}$ | $\begin{array}{r} \text { Field } \\ 4 \\ \hline \end{array}$ | $\begin{gathered} \text { Field } \\ 5 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 6 \\ \hline \end{gathered}$ | $\begin{aligned} & \text { Record } \\ & \text { Type } \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| zA | Awr | L. $¢$-1 | b | NK | b | head |
| $b$ | $b$ | b | b | NR | NP |  |
| $\mathrm{NBr}_{1}$ | $\mathrm{INT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBr}_{3}$ | $\mathrm{INT}_{3}$ |  |
| --- | --- | -- | --- | ${ }^{\text {NBT }}$ NR | ${ }^{\text {INT }}$ NR |  |
| $\mathrm{E}_{1}$ | $\mathrm{Y}\left(\mathrm{E}_{1}\right)$ | $\mathrm{E}_{2}$ | $\mathrm{Y}\left(\mathrm{E}_{2}\right)$ | $\mathrm{E}_{3}$ | $\mathrm{Y}\left(\mathrm{E}_{3}\right)$ |  |
| --- | --- | --- | --- | - | --- |  |
| --- | --- | --- | --- | $\mathrm{E}_{\mathrm{NP}}$ | $\mathrm{Y}\left(\mathrm{E}_{\mathrm{NP}}\right)$ | $\operatorname{tabl}^{\text {a }}$ |
| $\mathrm{EG}_{1}$ | $\mathrm{Es}_{1}$ | b | LF | NR | nP |  |
| $\mathrm{NBT}_{1}$ | $\mathrm{INT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| --- | --- | --. | --- | ${ }^{\mathrm{NBT}} \mathrm{NR}$ | ${ }^{\text {INT }}$ NR |  |
| $\mathrm{E}_{1}$ | $y_{1}\left(E_{1}\right)$ | $\mathrm{E}_{2}$ | $\mathrm{y}_{1}\left(\mathrm{E}_{2}\right)$ | $\mathrm{E}_{3}$ | $y_{1}\left(E_{3}\right)$ |  |
| --- | --- | --- | --- | -- | --- |  |
| -- | --- | --- | --- | $\mathrm{E}_{\mathrm{NP}}$ | $\mathrm{y}_{1}\left(\mathrm{E}_{\mathrm{NP}}\right)$ | tab1 |
| $\mathrm{EG}_{2}$ | $\mathrm{ES}_{2}$ | b | ${ }_{\text {LF }}$ | NR | NP |  |
| $\mathrm{NBT}_{1}$ | ${ }^{\text {inT }} 1$ | $\mathrm{NBT}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| --- | --- | --- | --- | ${ }^{\text {NBT }}{ }_{\text {NR }}$ | $\mathrm{INT}_{\mathrm{NR}}$ |  |
| $\mathrm{E}_{1}$ | $\mathrm{y}_{2}\left(\mathrm{E}_{1}\right)$ | $\mathrm{E}_{2}$ | $\mathrm{y}_{2}\left(\mathrm{E}_{2}\right)$ | $\mathrm{E}_{3}$ | $\mathrm{y}_{2}\left(\mathrm{E}_{3}\right)$ |  |
| --- | --- | -- | -- | -- | --- |  |
| --- | --- | --- | --- | $\mathrm{E}_{\mathrm{NP}}$ | $\mathrm{y}_{2}\left(\mathrm{E}_{\mathrm{NP}}\right)$ | TAB1 |
| =- | $=$ | =- | =-- | $\cdots$ | =- |  |
| ${ }^{E G}{ }_{\text {NK }}$ | $\mathrm{ES}_{\text {NK }}$ | b | ${ }_{\text {LF }}$ | NR | NP |  |
| $\mathrm{NBT}_{1}$ | $\mathrm{INT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| -- | --- | --- | -- | ${ }^{\mathrm{NBT}} \mathrm{NR}$ | ${ }^{1 N T}{ }_{\text {NR }}$ |  |
| $\mathrm{E}_{1}$ | $\mathrm{y}_{\mathrm{NK}}\left(\mathrm{E}_{1}\right)$ | $\mathrm{E}_{2}$ | $\mathrm{y}_{\mathrm{NK}}\left(\mathrm{E}_{2}\right)$ | $E_{3}$ | $\mathrm{y}_{\mathrm{NK}}\left(\mathrm{E}_{3}\right)$ |  |
| --- | --- | --- | -- | --- | --- |  |
| --- | --- | --- | --- | ${ }_{\text {E }}^{\text {NP }}$ | $\mathrm{y}_{\mathrm{NK}}{ }^{\left(E_{\mathrm{NP}}\right)}$ | ${ }^{\text {TAB1 }}$ |
| b | b | b | b | b | b | serd |

Field 7 (Cols. 67 through 70) $=$ MAT
Field 8 (Cols. 71 through 72) $=\mathrm{MF}$.
Field 9 (Cols. 73 through 75) $=$ MF (except for SEND card, where MT -0 )
Field 10 (Cols. 76 through 80) - card sequence number
This record is omitted if $\mathrm{NK}=1$ (only one subsection)


FILE 12: Multiplicities and Transition Probability Arrays

## $L \phi=1$ (Option 1, Multiplicities)

$$
y_{k}(E)=\frac{\sigma_{k}^{\gamma}(E)}{\sigma(E)}
$$

$\sigma(E)$ is given in File 3 for same MT number.

$$
\begin{aligned}
& y_{k}\left(E_{\gamma} \leftarrow E\right)=y_{k}(E) f_{k}\left(E_{\gamma} \leftarrow E\right) \\
& \int_{0}^{E_{\gamma}^{\max }} f_{k}\left(E_{\gamma}+E\right) d E_{\gamma}=1.000 \\
& Y(E)=\sum_{k=1}^{N K} y_{k}(E) \cdot
\end{aligned}
$$

bCD CARD-IMAGE FORMAT FOR FILE 12

## multiplicities and transition probabllity array

| $\begin{gathered} \text { Field } \\ 1 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 2 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Feld } \\ 3 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 4 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 5 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 6 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Record } \\ \text { Type } \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| zA | AWR | L¢ ${ }^{\text {2 }}$ | LG | ns | b | head |
| $\mathrm{Es}_{\text {NS }}$ | b | $b$ | b | (LG+1)*NT | NT |  |
| $\mathrm{Es}_{1}$ | $\mathrm{TP}_{1}$ | ${ }_{\left(G P_{1}\right)}$ | $\mathrm{Es}_{2}$ | $\mathrm{TP}_{2}$ | ${ }_{\left(G P_{2}\right)}$ |  |
| --- | -- | - | $\overline{\text { Es }}$ | ${ }^{\text {--- }}$ | ${ }_{\left(G \mathrm{P}_{3 \mathrm{H}}\right)}^{\text {-- }}$ |  |
| -- | --- | - | $\mathrm{ES}_{\mathrm{NT}}$ | ${ }^{\text {TP }}$ NT | ${ }^{\left(6 P_{N T}\right)}$ | LIST |
| b | b | b | b | b | b | send |


| 12024.0 | $2.4103 \mathrm{E}+11$ | 2 |  | 1 | 14 |  | 500212 | 64 | 476 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| .0000E+06 |  |  |  |  | , |  | 550212 |  | 47 |  |
| $4.2320 \mathrm{~F}+06$ | 0.045 | 4.1290E+06 | 0.045 |  | 1.30н0E. 06 | 0.91 | 50212 500212 |  | 4 |  |

F1eld 7 (Co1a. 67 through 70) $=$ MAT
F1eld 8 (cola. 71 through 72) $=$ MF


FILE 12: Multiplicities and Transition Probability Arrays

## $L \phi=2$ (Option 2, Transition Probability Arrays)

$A_{i}=\left(T P_{i}\right)\left(G P_{i}\right)$
$\frac{d \sigma}{d E_{\gamma}}\left(E_{\gamma}, E\right)=\sum_{j=1}^{N} \sum_{i=0}^{j-1} \sum_{m_{0}=j}^{N} \delta\left(E_{\gamma}-\varepsilon_{j}+\varepsilon_{i}\right) A_{j, i} \sigma_{m_{0}}(E) \overbrace{\ell=1}^{m_{0}-j} \sum_{m_{\ell}=j}^{m_{\ell-1}-1} T P_{m_{\ell-1}}, m_{\ell}$
$\sum_{i=0}^{N S-1} T P_{i}=1.0000$.

FILE 13: Photon Production Cross Sections
NK the total number of subsections for a reaction type section (MT number); i.e., the number of discrete photons, plus one if a photon continuum is given.

NR the number of interpolation ranges given.
NP the total number of neutron energy (E) points in the TABI tabulation to follow; i.e., the number of energy-cross-section pairs.
$\underline{N B T_{i}, I N T_{i}}$ the interpolation scheme for the ith interpolation region.
$E_{i}$ the neutron energy for the ith pair in the tabulation.
$\sigma_{\text {TOT }}^{\gamma}\left(E_{i}\right)$ the total photon production cross section for all subsections (not given if $\mathrm{NK}=1$ ),

$$
\sigma_{\text {TOT }}^{\gamma}\left(E_{i}\right)=\sum_{k=1}^{N K} \sigma_{k}^{\gamma}\left(E_{i}\right)
$$

$\sigma_{k}^{\gamma}\left(E_{i}\right)$ the partial photon production cross section for the kth subsection within a reaction type section; i.e., for a particular discrete photon or a photon continuum. In the case of a continuum, it is the integrated (over photon energy) cross section; the normalized photon energy distribution is given in File 15.
$\mathrm{EG}_{\mathrm{k}}$ the photon energy. For a continuous photon energy distribution, $E G_{k} \equiv 0.0$ should be used.
$\mathrm{ES}_{\mathrm{k}}$ the energy of the level from which the photon originates. If the level is unknown or if a continuous photon spectrum is produced, then $E S_{k} \equiv 0.0$ should be used.

LF the photon energy distribution law number, which presently has only two values defined.

LF $=1$, a normalized tabulated function (in File 15), and
$=2$, a discrete photon energy.


Field 7 (Cols. 67 through 70) $=$ MAT
Field 8 (Co1s. 71 through 72$)=$ MF $=13$
Field 9 (Cols. 73 through 75$)=$ MT (except for SEND card, where $\mathrm{MT}-0)$
ield 10 (Cols. 76 through 80 ) - card sequence number

Thic record omitted if NK = 1 (only one eubeaction)

FILE 13: Photon Production Cross Sections

$$
\begin{aligned}
\frac{d \sigma_{k}^{\gamma}}{d E_{\gamma}}\left(E_{\gamma}+E\right) & =\sigma_{k}^{\gamma}(E) f_{k}\left(E_{\gamma}+E\right) \\
\int_{0}^{E_{\gamma}^{\max }} f_{k}\left(E_{\gamma}+E\right) d E_{\gamma} & =1.000 \\
\sigma_{T O T}^{\gamma}(E) & =\sum_{k=1}^{N K} \sigma_{k}^{\gamma}(E)
\end{aligned}
$$

## bcd Card-image format for file 14

1.2024 E 042.4103 E 01

1
500214516680
photon angular distributions

| $\underset{1}{\text { Field }}$ | $\begin{gathered} \text { Field } \\ 2 \end{gathered}$ | $\begin{gathered} \text { Field } \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 6 \\ \hline \end{gathered}$ | $\begin{aligned} & \text { Record } \\ & \text { Type } \\ & \hline \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| zA | ALR | LIT1 | b | NK | 6 | head |
| b | b | b | b | b | b | SEND |

[^5]FILE 14: Photon Angular Distributions
LI =1 (Isotropic Distributions for all Photons)

$$
\begin{aligned}
\mathrm{P}_{\mathrm{k}}(\mu, \mathrm{E}) & \equiv 0.5 \\
\mathrm{a}_{\ell}^{\mathrm{k}}(\mathrm{E}) & \equiv 0.0, \quad \ell=1(1) \mathrm{NL}
\end{aligned}
$$

FILE 14: Photon Angular Distributions

## LI $=0$ (Anisotropic Distributions for Some Photons)

## LTT $=1$ (Legendre Coefficient Representation

NK the total number of subsections for a reaction type section (MT number); 1.e., the number of discrete photons, plus one if a photon continuum is given.

NI the number of isotropic photon angular distributions given in a section (MT number).
$E G_{k}$ the photon energy. For a continuous photon energy distribution, $\mathrm{EG}_{\mathrm{k}} \equiv 0.0$ should be used.
$\mathrm{ES}_{\mathrm{k}}$ the energy of the level from which the photon originates. If the level is unknown or if a continuous photon spectrum is produced, then $E S_{k} \equiv 0.0$ should be used.

NR the number of interpolation ranges given.
NE the number of neutron energy points given in a TAB2 record.
$\mathrm{NBT}_{i}, \mathrm{INT}_{\mathbf{i}}$ the interpolation scheme for the ith interpolation region.
Ei the neutron energy for the ith LIST record in a subsection.
NL $_{i}$ the number of Legendre coefficients used to represent the angular distribution at neutron energy $E_{1}$, excluding $a_{0}^{k}(E) \equiv 1.0$ that is understood; that is, $\ell=1$ (1) NL.
$a_{l}^{k}\left(E_{i}\right)$ the $\ell$ th Legendre coefficient for the angular distribution of the $k t h$ photon at incident neutron energy $E_{i}$
scd card-tmace format for flle 14
photon ancular distributions


| $\begin{gathered} \text { Frold } \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ \substack{ \\ \hline} \\ \hline \end{gathered}$ | $\begin{gathered} \begin{array}{c} \text { Held } \end{array} \\ \hline \end{gathered}$ | $\underset{\substack{\text { Preld } \\ 4}}{4}$ | $\underset{5}{\text { Piold }}$ | $\begin{aligned} & \text { Field } \\ & \hline 6 \end{aligned}$ | Record Type |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| za | Akr | LI-0 | 17T-1 | кк | ni | Head |
| ${ }^{\text {eg }}$ | $\mathrm{Es}_{1}$ | b | b | b | b | Conr |
| $\mathrm{EG}_{2}$ | $\mathrm{Es}_{2}$ | - | b | b | b | Conr |
| $=$ | $=$ | =-- | $=$ | =- | =- |  |
| ${ }^{\mathrm{EG}_{\mathrm{NI}}}$ | $\mathrm{Es}_{\mathrm{NI}}$ | b | b | b | b | conr |
| ${ }^{\text {E6 }} \mathrm{NI}+1$ | $\mathrm{Es}_{\mathrm{NI}+1}$ | b | b | NR | NE |  |
| $\mathrm{NBT}_{1}$ | $\mathrm{ImT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{ins}_{2}$ | ${ }^{\mathrm{NBr}_{3}}$ | $\mathrm{inT}_{3}$ |  |
|  |  |  |  | ${ }^{\text {NBT }} \mathrm{NR}^{\text {a }}$ | $\mathrm{INT}_{\text {NR }}$ | TAB2 |
| $\mathrm{a}_{1}^{\mathrm{NI}+1}\left(\mathrm{E}_{1}\right)$ | $\stackrel{N}{2}_{\mathbb{N}_{2}^{I_{1}}{ }_{1}\left(E_{1}\right)}$ | - | $\stackrel{-}{\square}$ | $\stackrel{\mathrm{NL}_{1}}{-}$ |  | List |
| $b$ | $\mathrm{E}_{2}$ | b | $\bigcirc$ | $\mathrm{NL}_{2}$ | b |  |
| $\mathrm{a}_{1}^{\mathrm{NT}+1}\left(\mathrm{E}_{2}\right)$ | $\mathrm{A}_{2}^{\mathrm{NT+1}\left(\mathrm{E}_{2}\right)}$ | --- | --- | --- |  | List |
| $=$ | $=$ | $=$ | $=$ | $=$ | $=$ |  |
| b | $\mathrm{E}_{\mathrm{NE}}$ | b | - | ${ }^{\text {N2 }}$ | b |  |
| ${ }_{1}^{a_{1}{ }^{\mathrm{NI}+1}\left(E_{\text {NE }}\right)}$ | $a_{2}^{\mathrm{NI}+1}\left(\mathrm{E}_{\mathrm{NE}}\right)$ | -- | --- | --- | ${ }_{8}^{\text {NLI }}$ N1 $\left(\mathrm{E}_{\mathrm{NSE}}\right)$ | Lis |
| ${ }^{\text {EG }}{ }_{\text {NI }+2}$ | $\mathrm{Es}_{\mathrm{NI}+2}$ | b | b | NR | ne |  |
| $\mathrm{NBT}_{1}$ | $\mathrm{INT}_{1}$ | $\mathrm{NBr}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{Nar}_{3}$ | $\mathrm{wnT}_{3}$ |  |
| -- | --- | - |  | ${ }^{\text {NET }}$ NR | $\mathrm{INT}_{\mathrm{NR}}$ | tab2 |
| ${ }^{\text {b }}$ | $\varepsilon_{1}$ | b | b | $\mathrm{NL}_{1}$ | ${ }^{\text {b }}$ |  |
| $a_{1}^{\mathrm{M}+2}\left(\mathrm{E}_{1}\right)$ | ${ }_{2}^{\mathrm{A}_{2}+2}\left(\mathrm{E}_{1}\right)$ | --- | --- | - | ${ }_{8 \mathrm{NL}}^{\mathrm{NI}+2}\left(\mathrm{E}_{1}\right)$ | List |
|  | $\mathrm{E}_{2}$ | b | b | $\mathrm{NL}_{2}$ | ${ }^{\text {b }}$ |  |
| ${ }_{1}^{(1+2}\left(E_{2}\right)$ |  | -- | -- | -- |  | $\underline{\text { isst }}$ |
| $=$ | =- | =- | $=$ | =- | =-- |  |
| ${ }^{\text {b }}$ | $\mathrm{E}_{\mathrm{NE}}$ |  | b | $\mathrm{Nr}_{\mathrm{NE}}$ | b |  |
|  | ${ }_{2}^{\mathrm{NI}+2}\left(\mathrm{E}_{\text {NE }}\right)$ | --- | --- | --- | $2_{\text {NL }}{ }^{\mathrm{NN}+2}\left(\mathrm{E}_{\mathrm{NS}}\right)$ | $\stackrel{\text { Lisr }}{ }$ |
| =- | --- | $=$ | $=$ | =-- | $=$ |  |
| $\mathrm{Ec}_{\text {NK }}$ | $\mathrm{Es}_{\mathrm{NK}}$ | b | b | мR | NE |  |
| ${ }^{\text {NBr }}{ }_{1}$ | $\mathrm{INT}_{1}$ | ${ }^{\mathrm{NBr}}{ }_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| --- | --- | -- | --- | ${ }^{\text {NET }}$ VR | ${ }^{\text {INT }}$ NR | ${ }^{\text {AB }}$ |
| ${ }^{\text {b }}$ | $\mathrm{E}_{1}$ | b | b | $\mathrm{NL}_{1}$ | b |  |
| $\mathrm{a}_{1}^{\mathrm{MK}}\left(\mathrm{E}_{1}\right)$ | $\mathrm{a}_{2}^{\mathrm{TK}}\left(\mathrm{E}_{1}\right)$ | --- | -- | -- | ${ }_{-\mathrm{NLL}^{\mathrm{NK}}\left(\mathrm{E}_{1}\right)}$ | LIST |
| $b$ | $E_{2}$ | b | b | $\mathrm{NL}_{2}$ | b |  |
| $\mathrm{ar}_{1}^{\mathrm{NK}}\left(\mathrm{E}_{2}\right)$ | $\mathrm{a}_{2}^{\mathrm{NK}}\left(\mathrm{E}_{2}\right)$ | --- | -- | -- | ${ }_{\mathrm{NLL}^{\mathrm{NK}}\left(\mathrm{E}_{2}\right)}$ | List |
| $=$ | $=$ | $=$ | $=$ | $=$ | = |  |
|  | $\mathrm{E}_{\text {sE }}$ | b | $b$ | $\mathrm{Nrase}_{\text {we }}$ | b |  |
|  | ${ }_{2}^{\mathrm{NK}}\left(\mathrm{E}_{\mathrm{NE}}\right)$ | --- | -- | --- | $2_{\text {MIL }}^{\text {NK }}$ ( ${ }_{\text {INE }}$ ) | List |
| b | b | b | b | b | b | SENT |



This page left blank because no data presently exist in ENDF format under the LTT = 1 option.

## $\underline{L I}=0$ (Anisotropic Distributions for Some Photons)

## $\underline{L T T}=1$ (Legendre Coefficient Representation)

$$
\begin{aligned}
p_{k}(\mu, E) & =\frac{2 \pi}{\sigma_{k}^{\gamma}(E)} \frac{d \sigma_{k}^{\gamma}}{d \Omega}(\Omega, E) \\
& =\sum_{\ell=0}^{N L} \frac{2 \ell+1}{2} a_{\ell}^{k}(E) p_{\ell}(\mu) \\
P_{\ell}(\mu) & =\frac{1}{2^{\ell} \ell!} \frac{d^{\ell}}{d \mu^{\ell}} \cdot\left(\mu^{2}-1\right)^{\ell} \\
a_{\ell}^{k}(E) & =\int_{-1}^{1} p_{k}(\mu, E) p_{\ell}(\mu) d \mu \\
\int_{-1}^{1} p_{k}(\mu, E) d \mu & =1.000 .
\end{aligned}
$$

FILE 14: Photon Angular Distributions

## $\underline{L I=0 \text { (Anisotropic Distributions for Some Photons) }}$

 LTT $=2$ (Tabulated Angular Distributions)NK the total number of subsections for a reaction type section (MT number); i.e., the number of discrete photons, plus one if a photon continuum is given.

NI the number of isotropic photon angular distributions given in a section (MT number).
$E G_{k}$ the photon energy. For a continuous photon energy distribution, $\mathrm{EG}_{\mathrm{k}} \equiv 0.0$ should be used.
$E S_{k}$ the energy of the level from which the photon originates. If the level is unknown or if a continuous photon spectrum is produced, then $E S_{k} \equiv 0.0$ should be used.
NR the number of interpolation ranges given.
NE the number of neutron energy points given in a TAB2 record.
$\mathrm{NBT}_{i}, \mathrm{INT}_{i}$ the interpolation scheme for the ith interpolation region.
$E_{i}$ the neutron energy for the ith TAB1 record in a subsection.
$\underline{\mu}$ the cosine of the reaction angle, $\mu=\cos \theta$.
$\mathrm{P}_{\mathrm{k}}\left(\mu, \mathrm{E}_{\mathrm{i}}\right)$ probability distribution function (normalized) at $\mu$ of photons produced by a neutron of incident energy $E_{i}$ for the $k t h$ photon energy.

| BCD CARD－IMAGE FOMGAT FOR FILE 14 photon angular distributions |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| $\underset{1}{\text { Putidid }}$ | $\xrightarrow{\text { qtald }}$ | $\stackrel{\text { rata }}{3}$ | $\stackrel{\text { Flo．ad }}{\substack{\text { c }}}$ | $\stackrel{\text { ratad }}{\substack{\text { rem }}}$ | ${ }_{\text {Fratd }}^{6}$ | Record |
|  | NMR | L1－0 | וтT＝2 | sk | ${ }^{\text {w }}$ | EhD |
| $\mathrm{Ec}_{1}$ | $\mathrm{ss}_{1}$ | － | － | b | b | com |
| $\mathrm{Ec}_{2}$ | $\mathrm{Es}_{2}$ | ¢ | b | － | 。 | conr |
| $=$ | $=$ | $=$ | $=$ | ＝－ | $=$ |  |
| ${ }_{\text {Ex }}^{\text {gr }}$ | $\mathrm{Es}_{\text {WI }}$ | b | 。 | － | b | cown |
|  |  | $\underset{\text { nss }}{\substack{\text { b }}}$ | $\underset{\substack{b \\ \mathrm{mur}_{2}}}{ }$ | $\underset{\text { wner }_{3}}{\mathrm{NR}^{2}}$ | $\begin{gathered} \mathrm{kI} \\ \mathrm{HKT}_{3} \end{gathered}$ |  |
|  |  |  |  |  | ${ }_{\text {InT }{ }_{\text {wer }}}$ | ma32 |
| b | $z_{1}$ | $\bigcirc$ | ${ }^{6}$ | nr | ${ }^{\text {nP }}$ |  |
| ${ }_{\text {sar }}^{1}$ | ${ }_{\text {wrr }}^{1}$ | ${ }^{335_{2}}$ | ${ }^{\text {min }}$ | $\underbrace{\substack{\text { ner }}}_{\text {ner }}$ | $\underbrace{\mathrm{LuT}_{3}}_{\substack{\text { unt }}}$ |  |
| $\overline{\mu_{1}}$ | （ $\overline{\left(u_{1}, R_{1}\right)}$ | $\nu_{2}$ | $\mathrm{P}_{\mathrm{MI}+1}\left(\overline{\mu_{2}} \overline{2}^{2} \mathrm{R}_{1}\right)$ | ${ }_{4}{ }_{4}$ | ${ }^{\text {P }}$ |  |
|  |  |  |  | ${ }_{4}{ }_{\text {xp }}$ |  | T381 |
| － | $\Sigma_{2}$ | b | b | к | ne |  |
| ${ }^{\text {net }}$ | ${ }_{\text {wr }}^{1}$ | $\mathrm{MHI}_{2}$ | ${ }^{1 \mathrm{mr}_{2}}$ | net $_{3}$ | $\mathrm{EwT}_{3}$ |  |
| $\mathrm{u}_{1}$ | $\left.\mathrm{P}_{\mathrm{M} 1+1} \mathrm{~L}_{1} \mathrm{~L}_{1}, \mathrm{R}_{2}\right)$ | $\mu_{2}$ | $p_{\text {PIT }}{ }^{\left(\mu_{2}, \mathrm{E}_{2}\right)}$ | ${ }_{\text {Nstrin }}$ |  |  |
| －－ | － | － | －－ | ${ }^{4} \mathrm{ne}$ |  | $\mathrm{TMS}^{\text {I }}$ |
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| $\stackrel{\mathrm{mm}_{1}}{-}$ | ${ }_{\text {LuT }}$ | $\stackrel{\mathrm{Nar}_{2}}{-}$ | $\stackrel{\text { rNT }}{ }$ |  | ${ }_{\substack{\text { miv }}}^{\mathrm{IvT}_{3}}$ |  |
| ${ }_{1}$ |  | ${ }_{2}$ |  | ${ }^{\mu}{ }_{3}$ | ${ }^{\text {P }}$ | Ten |
|  |  |  |  | $\sim_{\text {se }}$ | ${ }^{\text {P }}$ |  |
| 三－ | 三 | $\equiv$ | $\equiv$ | 三 | E |  |
| ${ }^{v_{\text {wrx }}}$ | ${ }_{\text {vsw }}$ | ${ }^{\circ}$ |  | NR | （1） |  |
| $\stackrel{\text { Bn7 }}{-}$ | $\stackrel{\text { int }}{--}$ | $\stackrel{\mathrm{mar}_{2}}{-}$ | $\stackrel{\text { LrT }}{-}$ |  |  | TMS2 |
|  |  | b | $\bigcirc$ | nR | np |  |
| ${ }_{\text {nar }}^{1}$ | ${ }_{\text {rut }}^{\text {rim }}$ | ${ }^{\mathrm{NBr}_{2}}$ | ${ }^{\text {nit }}$ | ${ }_{\text {NBT }}$ |  |  |
| ${ }_{1}$ | $P_{\text {PK }}\left({ }^{(1+1}+\varepsilon_{1} \varepsilon_{1}\right)$ | ${ }^{\mu}$ | $\mathrm{p}_{\text {MK }}{ }^{\left(\mathrm{L}_{2}, \mathrm{E}_{1}\right)}$ | ${ }^{4}$ |  | 1 |
|  |  |  |  | ${ }_{\text {wip }}$ | ${ }^{\text {max }}$ | 1 |
| $\mathrm{wrr}_{1}$ | ${ }_{\text {ma }}$ | $\mathrm{NBr}_{2}$ | ${ }^{\text {mTr }}$ | $\mathrm{ner}_{3}$ | ${ }^{1 \mathrm{wr}_{3}}$ |  |
| $\bar{\mu}_{1}$ |  |  |  | ${ }_{\text {Nis }}^{\substack{\text { ur }}}$ |  |  |
|  |  | － |  | ${ }^{\text {sis }}$ |  | Ts， 1 |
| ＝ | $=$ | $=$ | $=$ | ＝ | $=$ |  |
| b | ${ }_{\text {eng }}$ | － | － | NK | ne |  |
| ${ }_{\text {Nst }}^{1}$ | ${ }_{\text {mir }}$ | $\mathrm{NBT}_{2}$ | ${ }^{\text {inT }}$ |  | $\mathrm{int}_{3}^{\mathrm{INT}_{3}}$ |  |
| $\stackrel{\mu_{1}}{-}$ | ${ }^{\text {PMK }}$（ ${ }_{1}$ | $\stackrel{4}{2}$ |  | ${ }^{4}$ |  |  |
|  |  |  |  |  |  | saw |
|  |  |  |  |  | 。 |  |



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| 5．0000F－01 | 4．9083E－n］h．0n0ne－01 | ら．3174F－61 $7.001110 \mathrm{~F}-01$ | 5．71743－01150（1214 | 647447 |
| 00F | O499E－01 9．0100E－01 | 6．2209F－01 1．0000F OU1 | －．24332－01500214 |  |

## FILE 14: Photon Angular Distributions

## $L I=0$ (Anisotropic Distributions for Some Photons)

## LTT $=2$ (Tabulated Angular Distributions)

$$
\begin{aligned}
p_{k}(\mu, E) & \left.=\frac{2 \pi}{\sigma_{k}^{\gamma}(E)} \frac{d \sigma_{k}^{\gamma}}{d \underline{\Omega}} \underline{\Omega}, E\right) \\
\int_{-1}^{1} p_{k}(\mu, E) d \mu & =1.000 .
\end{aligned}
$$

FILE 15: Continuous Photon Energy Spectra
NC number of partial distributions used to represent a normalized energy distribution, $f\left(E_{\gamma} \nleftarrow E\right)$.
LF the photon energy distribution law number, which presently has only one continuous distribution defined: $L F=1$, a normalized tabulated function.

NR the number of interpolation ranges given.
NP the total number of neutron or photon energy points in the tabulation to follow.
$\mathrm{NBT}_{1}, \mathrm{INT}_{\mathbf{i}}$ the interpolation scheme for the ith interpolation range.
$E_{i}$ the neutron energy for the ith $T A B 1$ record.
$p_{1}\left(E_{i}\right)$ the probability given to the first partial distribution $\left[\mathrm{p}_{1}\left(\mathrm{E}_{\mathrm{i}}\right) \equiv 1.0\right.$ because only one distribution law is presently defined].

NE the number of neutron energy points given in a TAB2 record.
$E_{Y k}$ the photon energy for the $k$ th pair in a TAB1 record for the tabulated photon energy distribution.
$\underline{g\left(E_{\gamma k} \leftarrow E_{1}\right)}$ the probability distribution tabulated at photon energy $E_{\gamma_{k}}$ for an incident neutron energy $E_{1}$ (units of ev-1).
bcd card-image format for file 15
continuous photon energy spectra

| $\underset{1}{\text { Field }}$ | $\begin{gathered} \text { Field } \\ \underset{2}{ } \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 3 \end{gathered}$ | $\begin{gathered} \text { Field } \\ 4 \end{gathered}$ | $\begin{gathered} \text { Field } \\ 5 \end{gathered}$ | $\begin{gathered} \text { Field } \\ 6 \end{gathered}$ | Record <br> Type |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| zA | AWr | b | b | $\mathrm{Nc}^{\text {a }}$ | b | head |
| b | b | b | LF=1 | NR | NP |  |
| ${ }^{\mathrm{NBr}}{ }_{1}$ | $\mathrm{INT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| --- | --- | -- | --- | ${ }^{\text {NBT }}{ }_{\text {NR }}$ | $\mathrm{INT}_{\mathrm{NR}}$ |  |
| $\mathrm{E}_{1}$ | $\mathrm{P}_{1}\left(\mathrm{E}_{1}\right)$ | $\mathrm{E}_{2}$ | $\mathrm{P}_{1}\left(\mathrm{E}_{2}\right)$ | $\mathrm{E}_{3}$ | $\mathrm{P}_{1}\left(\mathrm{E}_{3}\right)$ |  |
| --- |  | --- | --- | ${ }_{\text {E }}^{\text {NP }}$ | $\mathrm{P}_{1}\left(\mathrm{E}_{\mathrm{NP}}\right)$ | TAB1 |
| b | b | b | b | NR | NE |  |
| $\mathrm{NBT}_{1}$ | $\mathrm{INT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| -- | --- | --- | --- | ${ }^{\text {NBT }}{ }_{\text {NR }}$ | $\mathrm{INT}_{\mathrm{NR}}$ | tAB2 |
| b | $\mathrm{E}_{1}$ | b | b | NR | nP |  |
| ${ }^{\text {NBT }}{ }_{1}$ | $\mathrm{INT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| --- | --- | --- | --- | ${ }^{\mathrm{NBT}} \mathrm{NR}^{\text {a }}$ | ${ }^{1 N T}{ }_{\text {NR }}$ |  |
| ${ }^{\varepsilon_{\gamma 2}}$ | $g\left(E_{\gamma 1}+\Sigma_{1}\right)$ | ${ }^{\mathrm{F}^{2}}$ | $g\left(E_{\gamma 2}+E_{1}\right)$ | ${ }_{E_{\gamma 3}}$ | $g\left(\mathrm{E}^{\mathbf{Y} 3}{ }^{\left.+\mathrm{E}_{1}\right)}\right.$ |  |
| -- | --- | --- | --- | $\mathrm{E}_{\gamma \mathrm{NP}}$ | $g\left(E_{\gamma \times \mathbb{P}}+E_{1}\right)$ | ${ }_{\text {TAB1 }}$ |
| b | $\mathrm{E}_{2}$ | b | b | NR | NP |  |
| ${ }^{\text {NBT }}{ }_{1}$ | $\mathrm{INT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{iNT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| -- | --- | --- | --- | ${ }^{\mathrm{NBT}} \mathrm{NR}^{\text {d }}$ | $\mathrm{INT}_{\mathrm{NR}}$ |  |
| ${ }^{\text {E }}{ }_{\gamma 1}$ | $g\left(E_{\gamma 1}+E_{2}\right)$ | ${ }^{\text {¢ }}{ }^{2}$ | $g\left(E_{\gamma 2}+E_{2}\right)$ | ${ }_{E_{\gamma 3}}$ | $8\left(\mathrm{E}_{\gamma 3}+\mathrm{E}_{2}\right)$ |  |
| -- | --- | --- | --- |  | $\mathrm{g}\left(\mathrm{E}_{\gamma \mathrm{NP}}+E_{2}\right)$ | TAB1 |
| =- | =- | $=$ | =- | = | =- |  |
| b | $\mathrm{E}_{\mathrm{NE}}$ | b | b | NR | NP |  |
| ${ }^{\text {NBT }} 1$ | $\mathrm{INT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| --- | -- | -- | --- | ${ }^{\text {NBT }}{ }_{\text {NR }}$ | $\mathrm{INT}_{\mathrm{NR}}$ |  |
| ${ }^{\text {E }}{ }_{1}$ | $8\left(\mathrm{E}_{\mathrm{Y} 1}+\mathrm{E}_{\mathrm{NE}}\right)$ | ${ }^{\mathrm{E}} \mathrm{Y}^{2}$ | $8\left(\mathrm{E}_{\gamma 2}+\mathrm{E}_{\mathrm{NE}}\right)$ | ${ }_{E}{ }_{\gamma}{ }^{3}$ | $\mathrm{g}^{\left(\mathrm{E}_{\mathrm{\gamma}}+{ }^{+\mathrm{E}_{\mathrm{NE}}}\right)}$ |  |
| -- |  | -- | - | ${ }_{\text {E }}{ }_{\text {NPP }}$ | $g\left(E_{Y N P}{ }^{+E_{N E}}\right)$ | TAB1 |
| b | b | b | b | b | b | SEND |

Field 7 (Cols. 67 through 70) $=$ MAT
Field 8 (CO1s. 71 through 72) $=\mathrm{MF}-15$
Field 9 (Cols. 73 through 75 ) MT (except for SEND card, where $M T=0$ )
Field 10 (Cols. 76 through 80) $=$ card sequence number
abecause only one law is presently defined, only NC -1 is possible.

$1.2024 \mathrm{E}+042.4103 \mathrm{E}+01 \quad 1$| 1 |
| :--- |
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|  | 2 | 2 |  |
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| 4.0 | E+06 $\quad 1.0$ | 2.0 | E+07 1.0 |
| 0.0 |  |  |  |
|  | 4.0 |  |  |

$$
{ }^{9.0000 E+06}
$$



$1.0000{ }^{12}+05$ $1.2000 \mathrm{E}+07$

$$
\begin{array}{llll}
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\end{array}
$$

$\qquad$

$$
\begin{aligned}
& 1.8556 E-08 \\
& 1.6000 \mathrm{E}+07
\end{aligned}
$$

5500215
500215
$\qquad$

$\qquad$




$\begin{array}{llllllll}9.5000 E & 06 & 3.1761 E-08 & 1.0500 E+07 & 2.4434 E-08 & 1.1500 E+07 & 1.7104 E-09500215 & 15 \\ 1\end{array} 133$ $1.2500 \mathrm{E}+07 \quad 1.0588 \mathrm{E}-08 \quad 1.3500 \mathrm{E}+07 \quad 4.0723 \mathrm{E}-0$
$\begin{array}{rlrl}19 & 2 & 19500215151135 \\ 1.0000 E+05 & 500215151136\end{array}$ $\begin{array}{lllll}1.0000 E+05 & 1.8057 E-26 & 2.5000 E+05 & 5.7203 E-09 & 7.5000 E+05 \\ 1.2500 E & 1.787 C E-09500215 & 15 & 1137\end{array}$




500215011
$\begin{aligned} & 6.5000 \mathrm{E}+0 \\ & 9.5000 \mathrm{E}+0\end{aligned}$

FILE 15: Continuous Photon Energy Spectra

$$
\begin{aligned}
& f\left(E_{\gamma}+E\right)=\sum_{j=1}^{N C} P_{j}(E) q_{j}\left(E_{\gamma}+E_{i}\right) \\
& \sum_{j=1}^{N C} p_{j}(E)=1.000
\end{aligned}
$$

$$
\int_{0}^{\mathbb{E}_{\gamma}^{\max }} f\left(E_{\gamma}+E\right) d E_{\gamma}=1.000
$$

$$
\int_{0}^{\mathrm{E}_{\gamma}^{\max }} \mathrm{q}_{j}\left(E_{\gamma}+E\right) d E_{\gamma}=1.000 .
$$

FILE 23: "Smooth" Cross Sections
NR the number of interpolation ranges given.
NP the total number of photon energy ( $E_{\gamma}$ ) points given in the tabulation to follow; i.e., the number of energy-cross-section pairs.
$\mathrm{NBT}_{i}, \mathrm{INT}_{i}$ the interpolation scheme for the ith interpolation range.
$E_{\gamma j}$ the photon energy for the $j$ th pair in the tabulation.
$\sigma\left(E_{\gamma j}\right)$ the photon interaction cross section.

| bcd Card-mage format for file 23 <br> "sмоотн" cross sections |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{gathered} \text { Field } \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 2 \end{gathered}$ | $\begin{gathered} \text { Field } \\ 3 \end{gathered}$ | $\begin{gathered} \text { Field } \\ 4 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ \hline \end{gathered}$ | $\begin{aligned} & \text { Field } \\ & \hline 6 \end{aligned}$ | $\begin{aligned} & \text { Record } \\ & \text { Type } \end{aligned}$ |
| za | AWR | b | b | b | b | head |
| b | b | b | b | NR | NP |  |
| $\mathrm{NBT}_{1}$ | $\mathrm{INT}_{1}$ | $\mathrm{NBT}_{2}$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| -- | -- | --- | -- | ${ }^{\mathrm{NBr}} \mathrm{NR}$ | ${ }^{1 N T}{ }_{\text {NR }}$ |  |
| $\mathrm{E}_{\boldsymbol{\gamma} 1}$ | $\sigma\left(E_{\gamma 1}\right)$ | ${ }^{\mathbf{\Sigma}}{ }^{2}$ | $\sigma\left(E_{\gamma 2}\right)$ | ${ }^{\text {E }}{ }_{\gamma}{ }^{3}$ |  |  |
| $\cdots$ | --- | -- | --- | - | - |  |
| -- | -- | - | - | ${ }_{\text {E }}^{\text {YNP }}$ | $\sigma\left(\mathrm{E}_{\gamma N \mathrm{P}}\right)$ | tabi |
| b | b | b | b | b | $\bigcirc$ | SEND |


| $5.00 \cdot 100+4$ | $1.17670+$ |  | 0 | 0 | 0 | 0 | 5023602 | 3 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $0.00000+0$ | 0.00000+ | 0 | 0 | 0 | 1 - | 49 | 5023502 | 284 |
| 49 |  |  |  |  |  |  | 5023602 | 285 |
| $1.00000+3$ | $1.59000+$ | 61.50000 | $36.43000+$ | 2.00000 | 33.21 |  | 5023602 | 286 |
| $3.00000+$ | 1.180004 | $53.92900+$ | $36.00000+$ | $43.92939+$ | 31.970 |  | 5023602 | 287 |
| $4.00000+$ | $1.84000+$ | $4.15600+$ | 31.66000 | 54.1564 | 32.28000 |  | 5023602 | 288 |
| $4.46500+$ | $1.90000+$ | $4.46545+$ | $32.19000+$ | 55.00000 | $31.65000+$ |  | 5023602 | 289 |
| $6.00 .1000+$ | 1.030004 | 5 8.00000+ | 3 4.81000 | 41.00000 | $42.64000+$ |  | 5023602 | 290 |
| $1.50000+$ | $8.69000+$ | 32.00000 | $43.91000+$ | 32.92000 | $41.37000+$ |  | 5023602 | 291 |
| $2.92029+$ | B. $51000+$ | $3.00000+$ | $47.87000+$ | 34.00000 | $43.70000+$ |  | 5023602 | 292 |
| $5.00000+$ | $2.01000+$ | $6.00000+$ | $41.22000+$ | 38.00000 | $45.43000+$ |  | 5023602 | 293 |
| $1.00000+$ | 2.870004 | $1.50000+$ | $58.97000+$ | 12.00000 | $53.92000+$ |  | 5023602 | 294 |
| $3.001500+$ | $1.25000+$ | $4.00000+$ | $55.71000+$ | $05.00000+$ | $53.21000+$ |  | 5023602 | 295 |
| $6.00000+$ | 2.06000t | $8.00000+$ | $51.06000+$ | 01.00000 | $68.63000-$ |  | 5023602 |  |
| $1.50000+$ | 3.05000- | $2.00000+$ | $61.87000-$ | 13.00000 | $69.97000-$ |  | 5023602 | 297 |
| $4.00000+$ | $6.65000-$ | $5.00000+$ | $64.98000-$ | 26.00000 | 6 3.89000- |  | 5023602 | 298 |
| 8.00000+ | 2.72000- | $21.00000+$ | 72.10000 | $1.50000+$ | $71.32000-$ |  | 5023602 | 299 |
| $2.00000+$ | 9.55000 | $3.00000+$ | 76.14000 | $4.00000+$ | 74.55000 |  | 5023602 | 3 CO |
| $5.00000+$ | 3.61000- | $6.00000+$ | 12.98000 | $3.00000+$ | $2.23000-$ |  | 5023602 | 3 Cl |
| $1.00000+8$ | 1.77000- | 3 |  |  |  |  | 5023602 | 3 C 2 |
|  |  |  |  |  |  |  |  |  |

Field 7 (Cols. 67 through 70) $=$ MAT


NR the number of interpolation ranges given.
NP the total number of recoil momentum (q) points given in the tabulation to follow; i.e., the number of momentum-form-factor pairs.
$\mathrm{NBT}_{1}, \mathrm{INT}_{i}$ the interpolation scheme for the ith interpolation range.
$\mathrm{q}_{\mathrm{i}}$ momentum of the recoil electron for incoherent scattering, or momentum of the atom for coherent scattering (in units of mos).
$H\left(q_{i} ; Z\right)$ is $K\left(q_{i} ; Z\right)$, the incoherent scattering function in the case of Compton scattering; or $G\left(q_{i} ; Z\right)$, the form factor in the case of coherent scattering.

| bCD Card mmage format for file 27 |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{gathered} \text { Field } \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 2 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ \hline \end{gathered}$ | $\begin{aligned} & \text { Field } \\ & \hline \end{aligned}$ | $\begin{gathered} \text { Field } \\ 5 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Field } \\ 6 \\ \hline \end{gathered}$ | Record <br> Type |
| zA | AWR | b | b | b | b | head |
| b | $z$ | b | b | NR | nP |  |
| $\mathrm{NBT}_{1}$ | $\mathrm{INT}_{1}$ | ${ }^{\text {NBT }} 2$ | $\mathrm{INT}_{2}$ | $\mathrm{NBT}_{3}$ | $\mathrm{INT}_{3}$ |  |
| -- | -- | -- | -- | ${ }^{\mathrm{NET}} \mathrm{NR}$ | $\mathrm{INT}_{\mathrm{NR}}$ |  |
| $\mathrm{q}_{1}$ | $\mathrm{H}\left(\mathrm{q}_{1} ; 2\right)$ | $\mathrm{q}_{2}$ | $\mathrm{H}_{(92} \mathbf{2}$ ) | $9_{3}$ | $\left.\mathrm{H}_{\left(\mathrm{q}_{3} ;\right.} \mathrm{z}\right)$ |  |
| - | --- | - | -- | --- | -- |  |
| --- | -- | --- | -- | ${ }^{9}{ }_{\text {MP }}$ | $\left.\mathrm{H}_{(9, \mathrm{NP}} ; 2\right)$ | tabl |
| b | b | b | b | b | b | SEND |

## FILE 27:

This page left blank because no data presently exist in ENDF File 27.

Field 7 (Cols. 67 through 70) $=\mathrm{MAT}$
Fteld 8 (Cols. 71 through 72) $=$ MF -27

teld 10 (Cols. 76 through 80 ) $=$ card sequence number

[^6]
[^0]:    *There should be no more than one energy continuum for each MT number used. If the decomposition of a continuum into several parts is desired, this can be accomplished in File 15.

[^1]:    *If the total number of discrete photons and photon continua is one (NK = 1), then this TABl record is omitted.

[^2]:    *There should be no more than one energy continuum for each MT number used. If the decomposition of a continuum into several parts is desired, this can be accomplished in File 15.

[^3]:    *If the total number of discrete photons and photon continua is one ( $\mathrm{NK}=1$ ), then this TAB1 record is omitted.

[^4]:    *Note that the subscript $k$ that occurs in the sections describing Files 12 and 13 has been dropped from $f\left(E_{\gamma} \leftarrow E\right)$. This is done because only one energy continuum is allowed for each MT number, and the subscript $k$ has no meaning in File 15. It is, in fact, the NKth subsection in File 12 or 13 that contains the production data for the continuum.

[^5]:    Field 7 (Cols. 67 through 70) $=$ MAT
    Field 8 (Cole 71 through 72) $=$ MF $=14$
    for SEND card, where MT - 0 )
    

[^6]:    FILE 27: Atomic Form Factors or Scattering Functions
    $E_{\gamma}^{\prime}$ is the scattered photon energy.
    $\mu=\cos \theta$
    $\begin{aligned} \alpha & =E_{\gamma} / m_{0} c^{2} \\ q & =\alpha\left[1+\left(\frac{\alpha^{\prime}}{\alpha}\right)^{2}-2 \mu\left(\frac{\alpha^{\prime}}{\alpha}\right)\right]^{\frac{1}{2}}, \text { incoherent scattering }\end{aligned}$
    $=\alpha[2(1-\mu) 2]^{\frac{1 / 2}{2}}$, coherent scattering.
    $\frac{\mathrm{d} \sigma_{i}}{\mathrm{~d} \mu}=2 \mathrm{~K}(\mathrm{q} ; 2) \frac{\mathrm{d} \sigma_{\mathrm{c}}}{\mathrm{d} \mu}$
    $\frac{{ }^{d \sigma}{ }_{c o h}}{d \mu}=\pi r_{0}^{2} z^{2}\left(1+\mu^{2}\right) G(q ; z)$.

